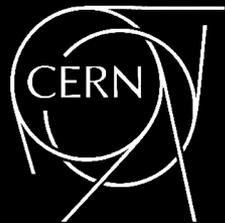


# ALPs at Future $e^-p$ Colliders

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WITS  
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ATLAS  
EXPERIMENT

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of the South African Institute of Physics

# Axion-Like Particles

- Axion-Like Particles (or ALPs) are **neutral pseudo scalar** pseudo-nambu Goldstone bosons.
- Interaction to SM particles arise from explicit **Global Peccie-Quinn  $U(1)_{PQ}$  symmetry**.
- Either shift invariant and/or anomalous couplings interactions  $\left\{ \begin{array}{l} \frac{\partial_\mu a}{f_a} \bar{\psi} \gamma^\mu \psi \\ \frac{a}{f_a} X_{\mu\nu} \tilde{X}^{\mu\nu} \end{array} \right.$
- ALPs interactions with SM particles have a derivative character: They grow with momentum

# Motivations for Axions and ALPs



# Axion-Like Particles

$$\mathcal{L}_{eff} \supset e^2 \frac{a}{f_a} g_{\gamma\gamma} F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{2e^2}{c_w s_w} \frac{a}{f_a} g_{\gamma Z} F_{\mu\nu} \tilde{Z}^{\mu\nu} + \frac{e^2}{s_w^2 c_w^2} \frac{a}{f_a} g_{ZZ} Z_{\mu\nu} \tilde{Z}^{\mu\nu} + \frac{e^2}{s_w^2} \frac{a}{f_a} g_{WW} W_{\mu\nu} \tilde{W}^{\mu\nu}$$

$f_a = 1 \text{ TeV}$  : New Physics Scale

- Classical Searches: ALPs coupling to gluons and photons

$WW, ZZ, \gamma Z$  and  $\gamma\gamma$

- ALPs couplings to EWK bosons:

$$g_{\gamma\gamma} = C_{WW} + C_{BB},$$

$$g_{Z\gamma} = c_w^2 C_{WW} - s_w^2 C_{BB},$$

$$g_{ZZ} = c_w^4 C_{WW} + s_w^4 C_{BB},$$

$$g_{WW} = C_{WW}.$$

- ALP-gauge interactions at ATLAS and CMS:

- Mono-X

- **New Idea: Non-resonant**

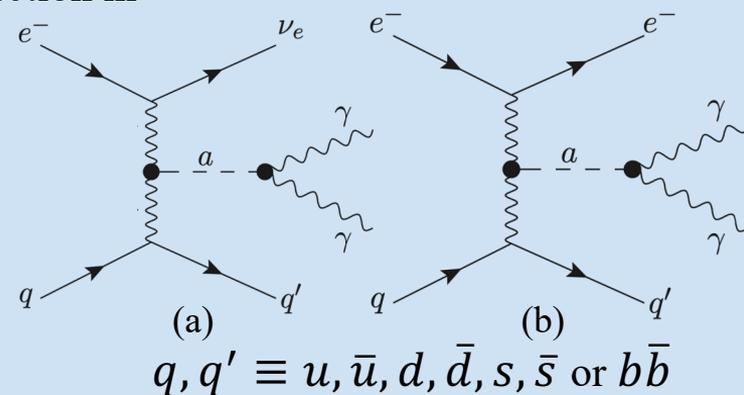
- Resonant

# Our Approach

- Leading order and matrix-element level single ALP production in

- Charged Current (CC) through (a)  $W^\pm$  - Fusion

- Neutral Current (NC) through (b)  $ZZ, \gamma Z$  and  $\gamma\gamma$ .



- Signals independent on ALP mass and width up to

$$m_a \ll 300 \text{ GeV}$$

# Our Approach

- The relevant decay widths are modelled as follows:

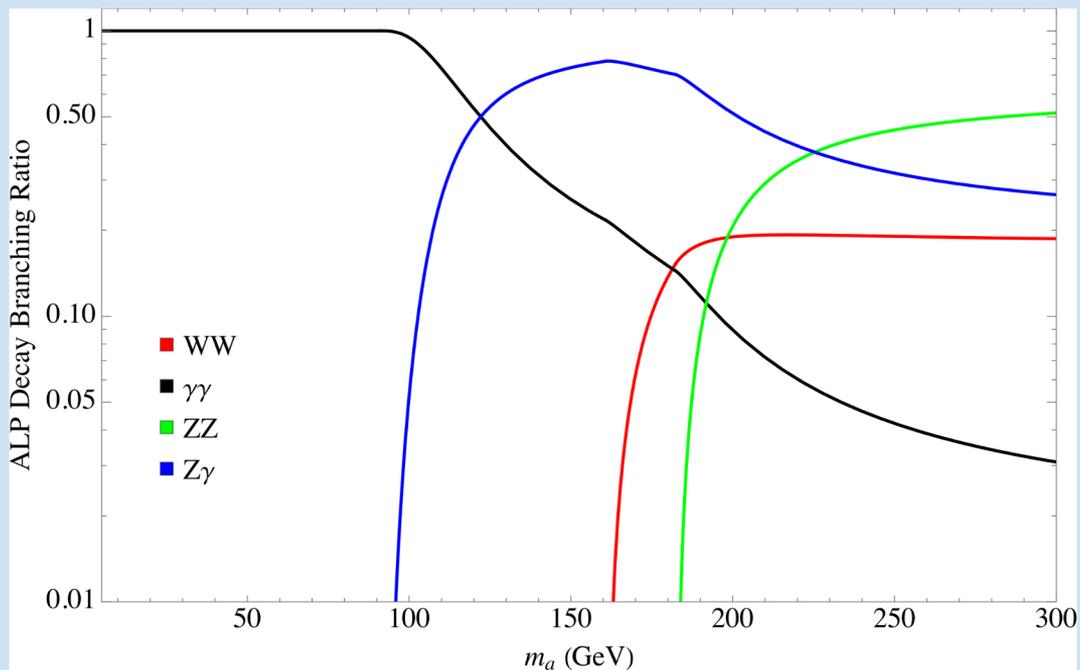
$$\Gamma(a \rightarrow \gamma\gamma) \equiv \Gamma_{\gamma\gamma} = \frac{e^4}{4\pi f_a^2} |g_{\gamma\gamma}|^2 m_a^3, \quad \Gamma(a \rightarrow \gamma Z) \equiv \Gamma_{\gamma Z} = \frac{e^4}{2\pi f_a^2 c_{\theta_W}^2 s_{\theta_W}^2} |g_{\gamma Z}|^2 m_a^3 \left(1 - \frac{m_Z^2}{m_a^2}\right)^3$$

$$\Gamma(a \rightarrow ZZ) \equiv \Gamma_{ZZ} = \frac{e^4}{4\pi f_a^2 c_{\theta_W}^4 s_{\theta_W}^4} |g_{ZZ}|^2 m_a^3 \left(1 - 4 \frac{m_Z^2}{m_a^2}\right)^{3/2},$$

$$\Gamma(a \rightarrow W^\pm) \equiv \Gamma_{W^\pm} = \frac{e^4}{8\pi f_a^2 s_{\theta_W}^4} |g_{WW}|^2 m_a^3 \left(1 - \frac{m_W^2}{m_a^2}\right)^{3/2}$$

**NOTE: We considered variable decay widths**

# Our Approach



- Branching ratio for the considered ALPs production channels.

# Our Approach

- The local significance and discovery limits are given as follows:

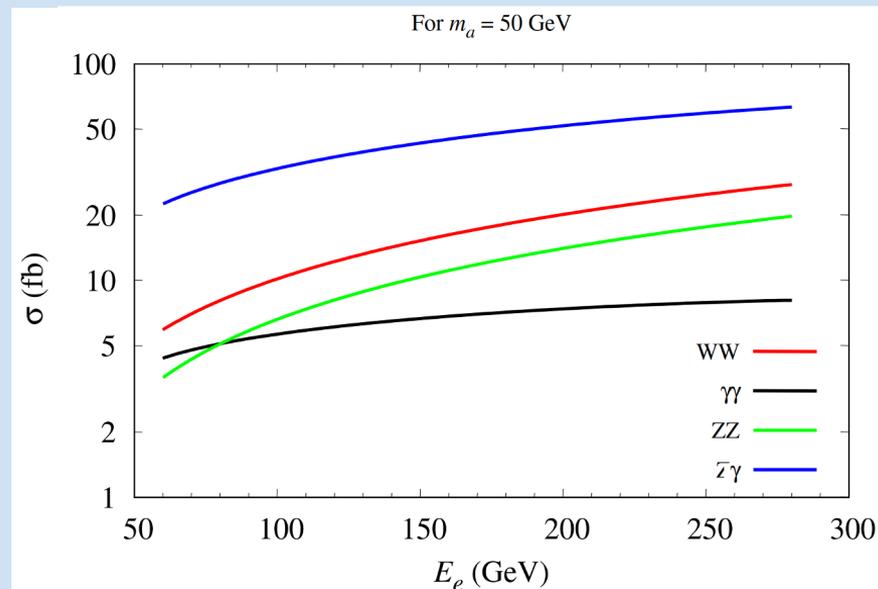
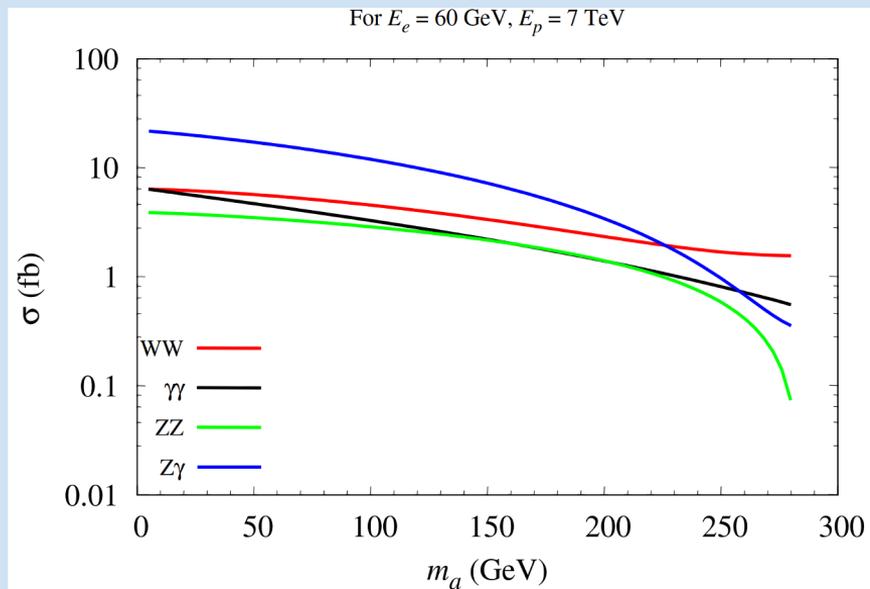
$$N_{SD} = \frac{S}{\sqrt{S + B + (\delta_S \cdot S)^2 + (\delta_S \cdot B)^2}}$$

- The  $\chi^2$  definition used to constraint ALP-gauge couplings  $g_{ij}$  is given as:

$$\chi^2 = \sum_{k=1}^n \left( \frac{N_k(g_{ij}) - N_k^{SM}}{\Delta N_k} \right)^2, \quad \Delta N_k = \sqrt{N_k^{SM} (1 + \delta_S^2 N_k^{SM})}$$

Note: 2 cases considered – (1) **one-bin  $\chi^2$  analysis** and (2) **multiple-bin  $\chi^2$  analysis**.

# Our Approach



- **$Z\gamma$**  Can not be separate at production level due to the interference of  **$\gamma\gamma$**

# ALP production in $e^-p$ collider

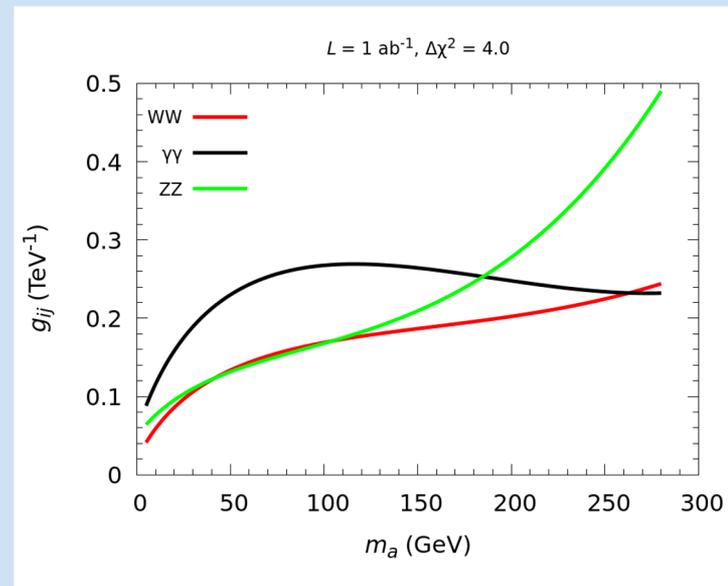
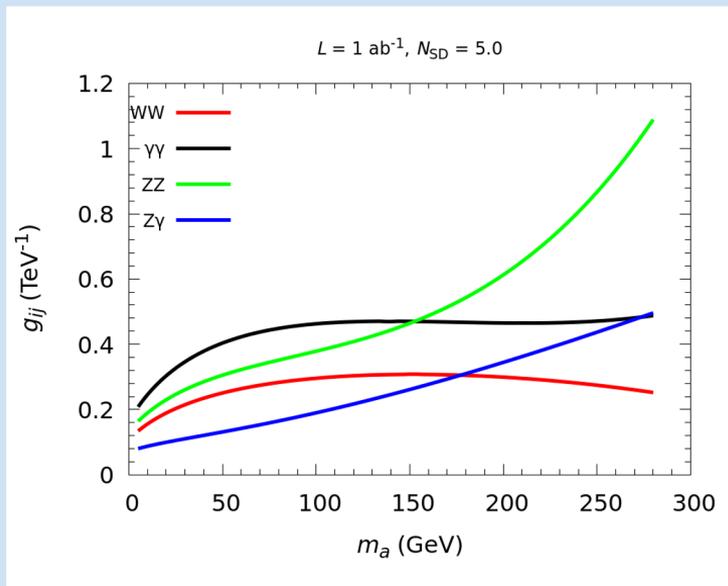
- ALP-gauge couplings are probed by direct production of ALP through VBF processes.
- We base our study on LHeC environment, employing **7 TeV** proton beam from LHC and electrons from Energy Recovery Linac (ERL):  $\sqrt{s} = \mathbf{1.3\ TeV}$ .

**FeynRules**  $\longrightarrow$  **MadGraph5\_aMC@NLO**  $\longrightarrow$  **Pythia-PGS**  $\longrightarrow$  **Delphes**

# Analysis, Observable and Results

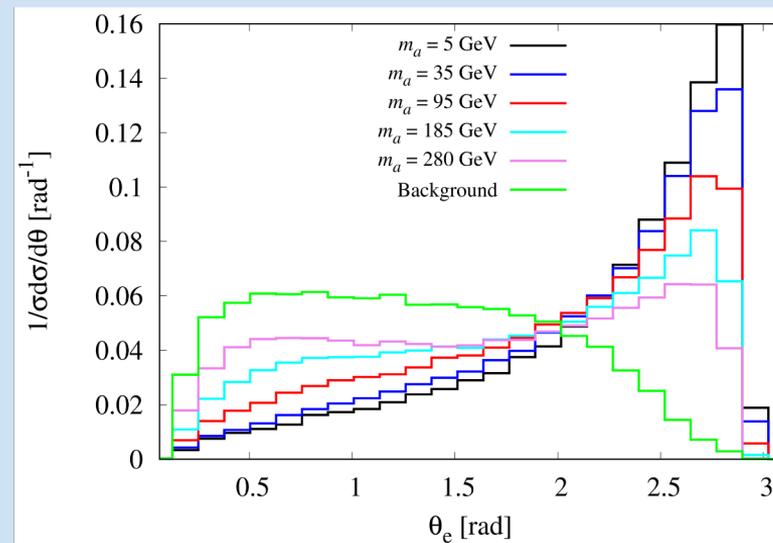
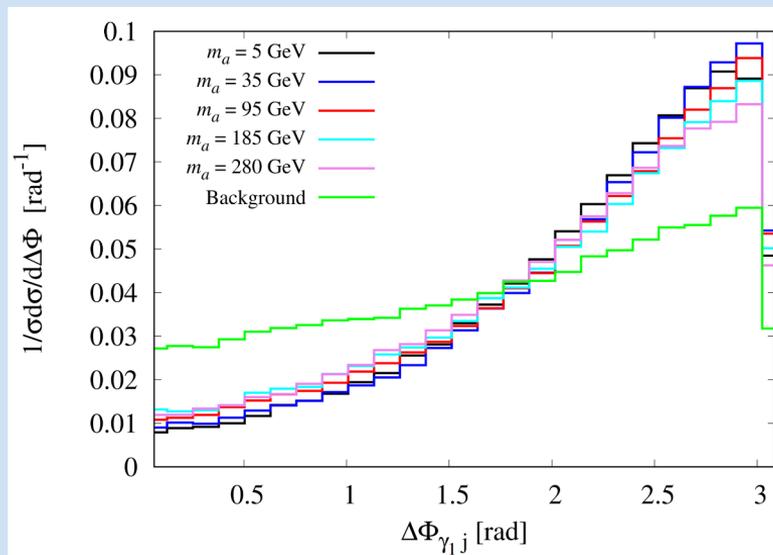
- To optimize signal events over leading SM backgrounds, additional cuts were applied as follows:
  - for all channels:  $p_T^{j,e^-, \gamma} > 20 \text{ GeV}$
  - **WW**:  $0 < |\eta_\gamma| < 3, 0 < |\eta_j| < 4$
  - **$\gamma\gamma$** :  $-2 < |\eta_\gamma| < 3, -2 < |\eta_j| < 5, -2.5 < |\eta_{e^-}| < 2$
  - **ZZ**:  $0 < |\eta_\gamma| < 3, 0 < |\eta_j| < 5, 0 < |\eta_{e^-}| < 5$
  - **$\gamma Z$** :  $0 < |\eta_\gamma| < 3, 0 < |\eta_j| < 5, -2.5 < |\eta_{e^-}| < 1$
- Analysis was performed at **5% allowed systematic error**,  **$1 \text{ ab}^{-1}$  luminosity** and **80% left polarized electron beam**

# Analysis, Observable and Results



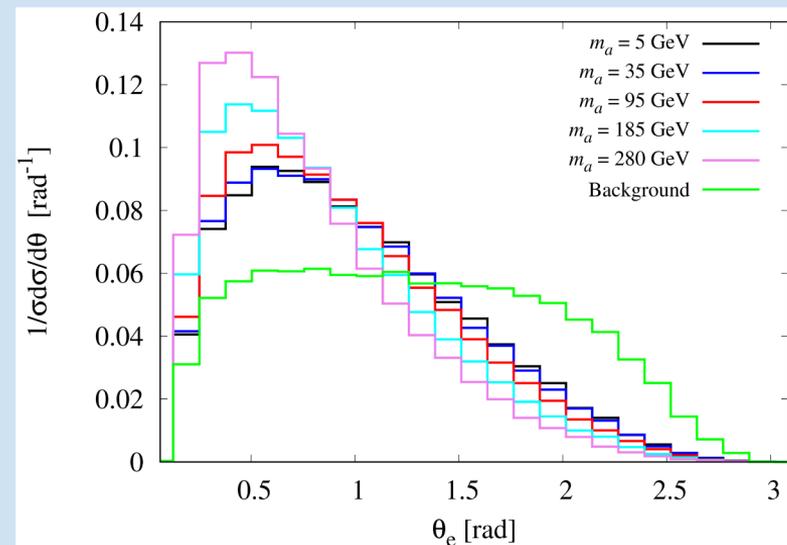
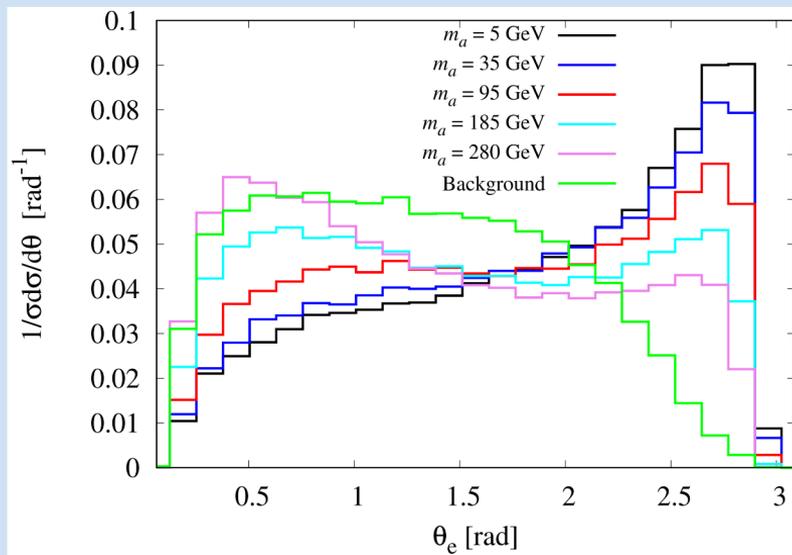
- Upper bounds on effective couplings as a function of Axion mass, based on the local significance at **5 $\sigma$**  discovery limit (**left**) and total differential cross section (**right**).

# Analysis, Observable and Results



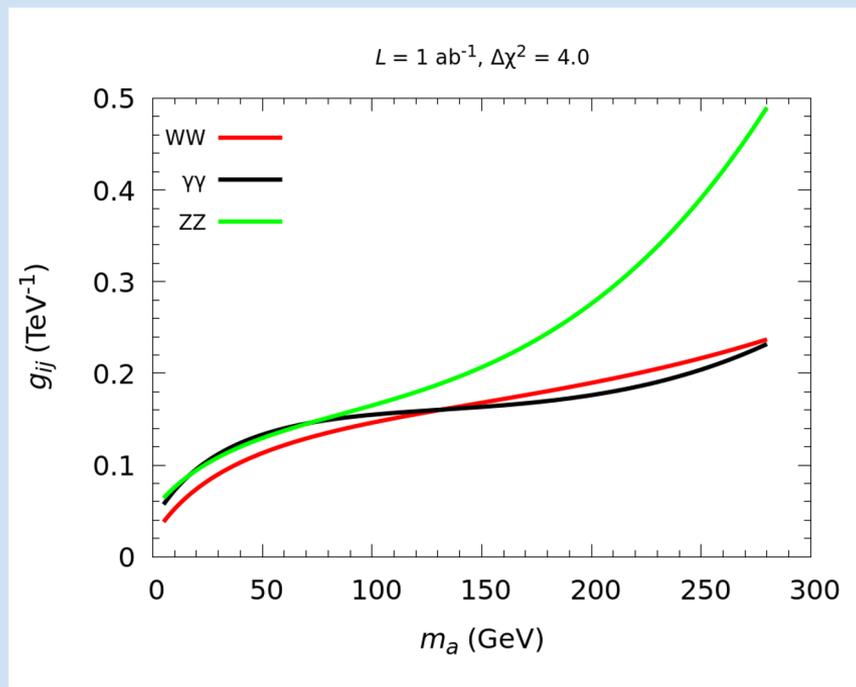
- Differential distributions for  $WW$ -fusion and  $\gamma\gamma$ -fusion, respectively.

# Analysis, Observable and Results



- Differential distributions for  $Z\gamma$  - fusion and  $ZZ$  - fusion, respectively.

# Analysis, Observable and Results



- Upper bounds on effective couplings,  $g_{ij}$  as a function of axion mass,  $m_a$  based on **sensitive differential distributions**

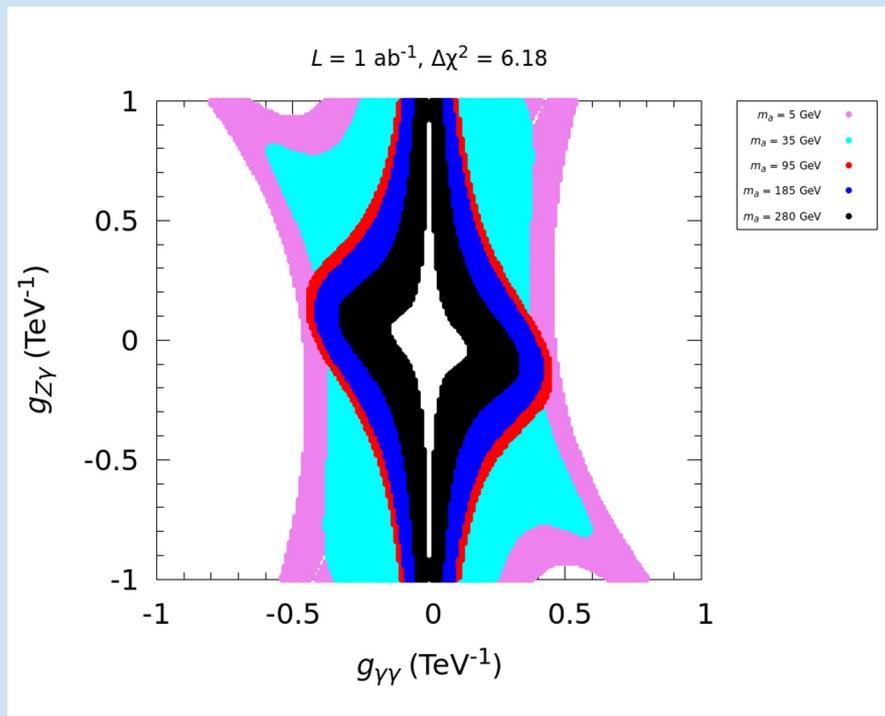
# Analysis, Observable and Results

- The  $g_{\gamma Z} = g_i$  and  $g_{\gamma\gamma}$  separation problem was solved by considering:

$$\sigma_{BSM} = g_{\gamma\gamma}^2 (g_i^2 \sigma_i^2 + g_i g_{\gamma\gamma} \sigma_{i\gamma\gamma} + g_{\gamma\gamma}^2 \sigma_{\gamma\gamma}) \times \mathcal{B}_{a \rightarrow \gamma\gamma}$$

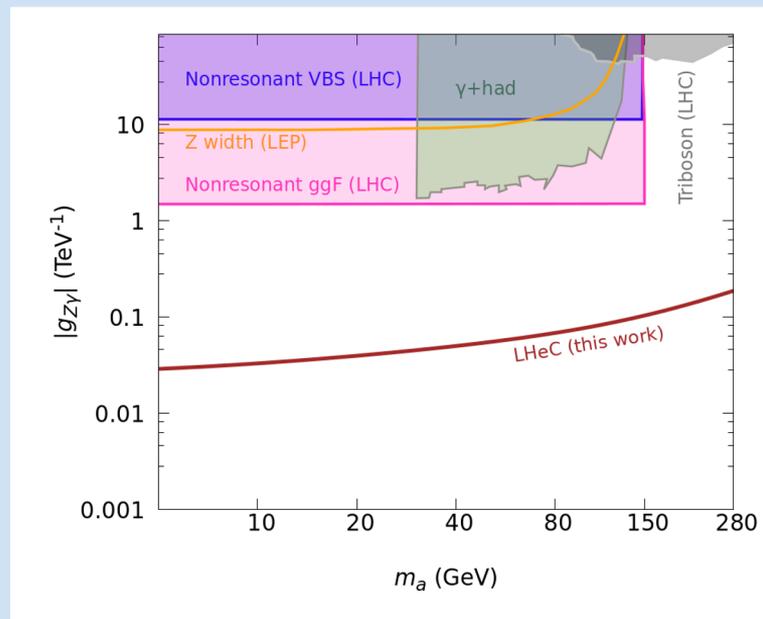
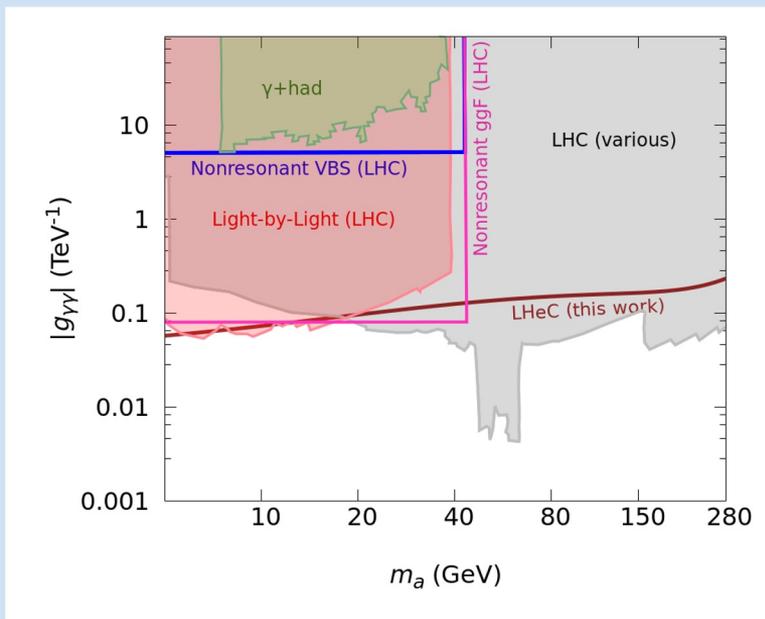
$g_{\gamma\gamma}$	$g_{\gamma Z} = g_i$
1	1
1	0
-1	1

# Analysis, Observable and Results



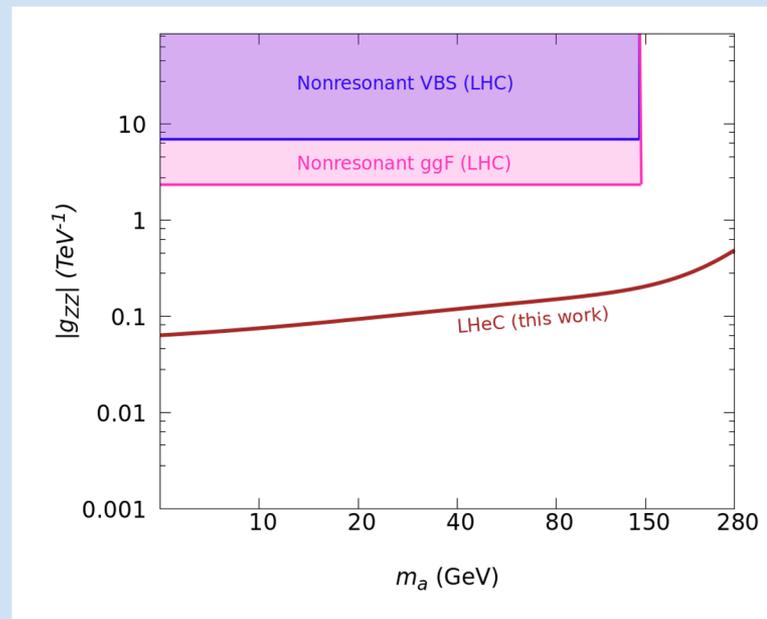
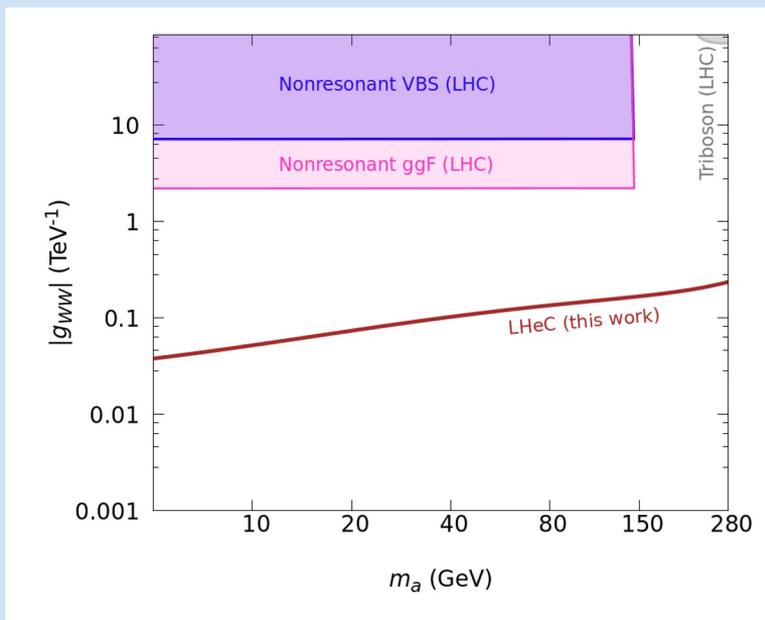
- 2-parameter upper bound with all **3 cases** taken into consideration due to the **SEPARATION PROBLEM**

# Experimental Comparison



- Various LHC, **Light-by-light scattering (LHC)**, **Nonresonant VBS**, **Nonresonant ggF (LHC)**,  **$\gamma$  + had (LHC)**, **LHeC (This work)**, **Triboson (LHC)** and **Z width (LEP)**.

# Experimental Comparison



- **Nonresonant VBS, Nonresonant ggF (LHC), Triboson (LHC)**

# Conclusion

- We conclude that an electron-proton collider running at high energies with a high luminosity would have great potential in searching for the ALPs.
- These limits serve as approximate predictions that can be investigated further if the electron energy ( $E_e$ ) is increased to higher values.
- The limits on  $g_{WW}$ ,  $g_{ZZ}$  and  $g_{YZ}$  comparing to available studies in different collider scenarios are better at LHeC for the considered range of  $m_a$ , whereby, the limits on  $g_{\gamma\gamma}$  are competitive with respect to few scenarios.
- VBF process used to probe ALPs provides a complementary window to the opportunities offered by other high energy collider experiments.
- Although, exploring high masses beyond 200 GeV is less likely due to the limited cross section achievable with the available energy and luminosity.



**Welcome to the "Future"**

**Thank You!**

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