

Weak Gravity Conjecture for dilaton de Sitter black holes in extra dimension

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Introduction

- ▶ The Weak Gravity Conjecture

In a flat space-time, for a single $U(1)$ gauge theory with gauge coupling g , there is at least one state that has a charge q bigger than its mass m , measured in Planck units

$$\frac{m}{M_P} < \sqrt{2}gq$$

(where M_P is the reduced Planck mass)

Arkani-Hamed, Motl, Nicolis, Vafa

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- ▶ For a first check we can look at the SM

$$|gq| \sim 0.3, \quad m_e = 0.511 \text{keV}, \quad M_P = 2 \times 10^{-18} \Rightarrow .42 > 2.5 \times 10^{-19}$$

Actually, all the fundamental particles of the SM verify this conjecture

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- ▶ One of the main motivations for the WGC is that it is a kinematic requirement that allows extremal BHs to decay.
- ✓ Since we are looking for charged objects, we consider the Reissner-Nordström metric:

$$ds^2 = \left(1 - \frac{2M}{r} + \frac{Q^2}{r^2}\right) dt^2 + \frac{dr^2}{1 - \frac{2M}{r} + \frac{Q^2}{r^2}} + r^2 d\Omega^2$$

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- ★ $M > Q$: This solution has two horizons, and describe a BH
- ★ $M = Q$: The two horizons are combined, (extremal BH)
- ★ $M < Q$: There is no horizon, the solution is a naked singularity and it is a non-physical solution.

- ▶ Let's consider an extremal BH, $M = Q$

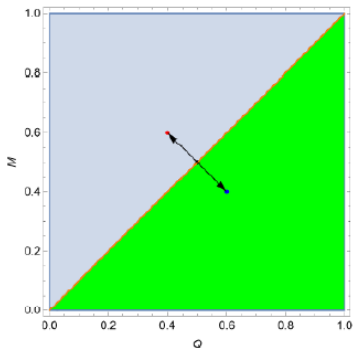
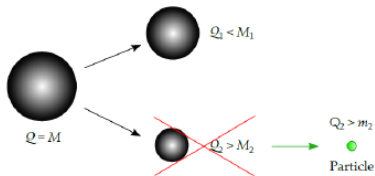


Figure: Extremal black holes can only decay into a particle that respects the WGC and a non-extremal black hole.

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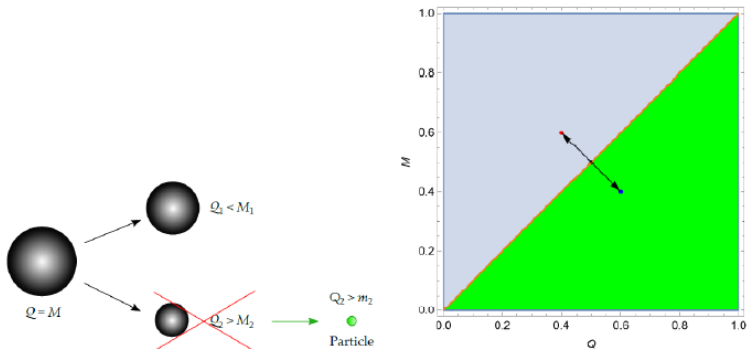


Figure: Extremal black holes can only decay into a particle that respects the WGC and a non-extremal black hole.

- ★ The WGC is thus seen to open up a possibility for extremal black holes to decay into non-extremal.

If we generalize to the WGC to the case with non-vanishing vacuum energy and also we add a dilaton field..

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- ▶ We will consider a dilaton coupling $e^{-2\alpha\phi} F^2$

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This leads to the Einstein-Maxwell-Dilaton action in $4D$:

$$S = \int d^4x \sqrt{-g} [R - \partial_\mu \phi \partial^\mu \phi - e^{2\alpha\phi} F^2 - V(\phi)]$$

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This action on n -dimension

$$S = \int d^n x \sqrt{-g} [R - \frac{4}{n-2} \partial_\mu \phi \partial^\mu \phi - e^{\frac{4\alpha\phi}{n-2}} F^2 - V(\phi)]$$

Strategy:

- ▶ 1) First we need to construct black hole solutions.
- ▶ 2) Next, we need to find the relation between mass and charge for extremal black holes.
- ▶ 3) Finally, we identify the WGC inequality corresponding to the existence of a super-extremal state.

$$\Lambda = 0 \text{ and } 4D$$

▶ $\alpha = 0$

$$g_{00} = \left(1 - \frac{r_+}{r}\right)\left(1 - \frac{r_-}{r}\right) \Rightarrow \text{R.N.BH}$$

- ★ No Naked singularity $\Rightarrow Q^2 < M^2 \Rightarrow r_{\pm} = M \pm \sqrt{M^2 - Q^2}$,
 $r = 0$ is a time- like singularity
 r_{\pm} are inner and event horizons

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▶ $\alpha \neq 0$

$$g_{00} = \left(1 - \frac{r_+}{r}\right)\left(1 - \frac{r_-}{r}\right)^{\frac{1-\alpha^2}{1+\alpha^2}}$$

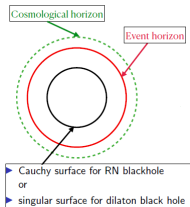
- ★ No Naked singularity $\Rightarrow Q^2 e^{2\alpha\phi_0} < (1 + \alpha^2)M^2$.
 r_- is a space- like singularity
 r_+ is the only event horizon of the black hole.
Gibbons-Maeda Garfinkle-Horowitz-Strominger

dS space : $\Lambda > 0$, n dimensions

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$$\begin{aligned}
 ds^2 = & - \left\{ \left[1 - \left(\frac{r_+}{r} \right)^{n-3} \right] \left[1 - \left(\frac{r_-}{r} \right)^{n-3} \right]^{1-\gamma(n-3)} - H^2 r^2 \left[1 - \left(\frac{r_-}{r} \right)^{n-3} \right]^\gamma \right\} dt^2 + \\
 & \left\{ \left[1 - \left(\frac{r_+}{r} \right)^{n-3} \right] \left[1 - \left(\frac{r_-}{r} \right)^{n-3} \right]^{1-\gamma(n-3)} - H^2 r^2 \left[1 - \left(\frac{r_-}{r} \right)^{n-3} \right]^\gamma \right\}^{-1} \\
 & \left[1 - \left(\frac{r_-}{r} \right)^{n-3} \right]^{-\gamma(n-4)} dr^2 + r^2 \left[1 - \left(\frac{r_-}{r} \right)^{n-3} \right]^\gamma d\Omega_{n-2}^2.
 \end{aligned} \tag{2.1}$$

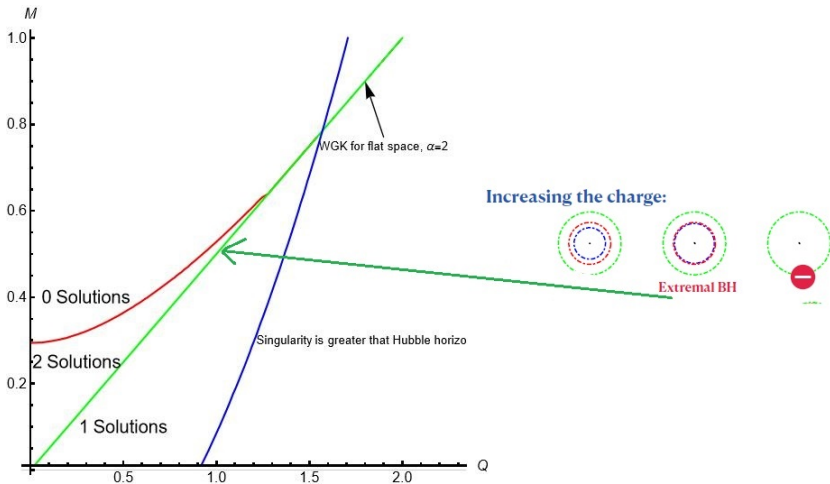
- ★ $H^2 = \frac{|\Lambda|}{3} \Rightarrow$ Hubble parameter. Gao-Zhang/ Elvang-Friedman-Liu
- ★ Note that r_- still indicates the coordinate of a singular surface,
- ★ r_+ is not an event horizon.
- ★ r_c is a cosmological horizon.



- The (0,0) component of the metric :

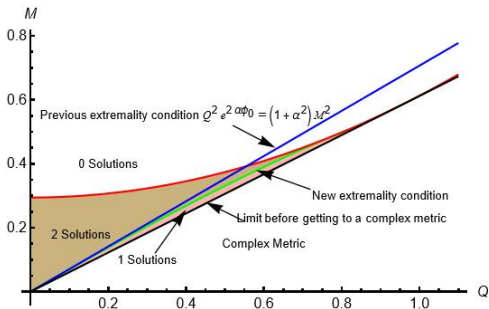
$$f(r) = \left[1 - \left(\frac{r_+}{r} \right)^{n-3} \right] - H^2 r^2 \left[1 - \left(\frac{r_-}{r} \right)^{n-3} \right] \alpha^2 - \frac{(n-3)^2}{n-1}$$
$$\alpha_c = \frac{(n-3)}{\sqrt{n-1}}$$

- ▶ $\Lambda > 0$ $\alpha = 2$ ($\alpha > \alpha_c = 1$) and $n = 5$

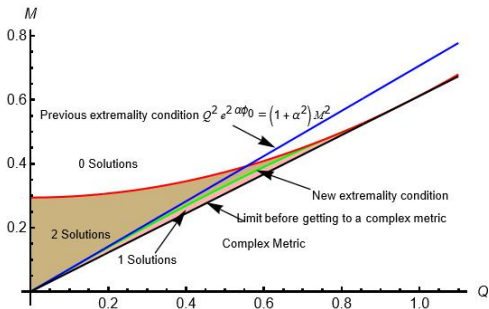


- ▶ WGC in flat space and 5 dimensions: $r_+ = r_- \Rightarrow M^2 = \frac{Q^2}{4}$

► $\Lambda > 0$ $\alpha = 1$ ($\alpha = \alpha_c = 1$) and $n = 5$

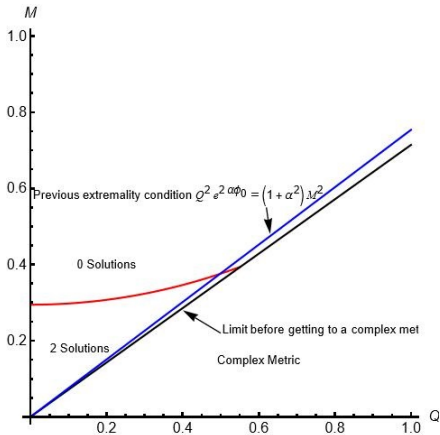


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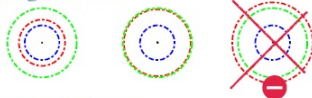


$$M^2 < \frac{Q^2}{2} + \frac{\sqrt{2}Q^2 H^2}{4\pi} - \frac{3Q^4 H^4}{16\pi^2} + \mathcal{O}(H^6)$$

▶ $\Lambda > 0$ $\alpha = 0.5$ and $n = 5$



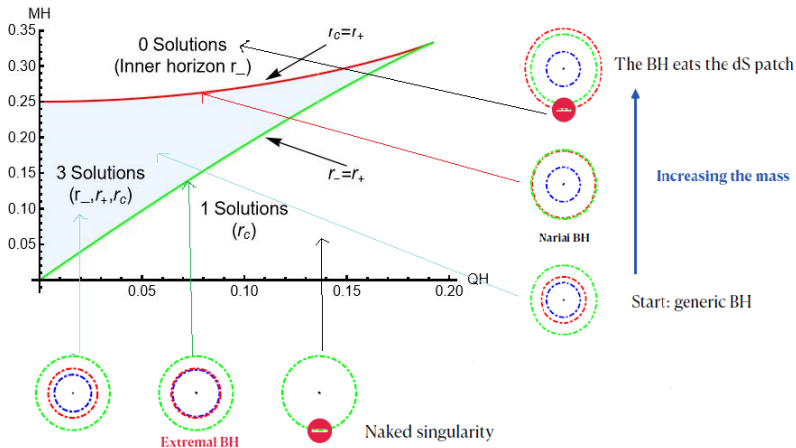
Increasing the mass:



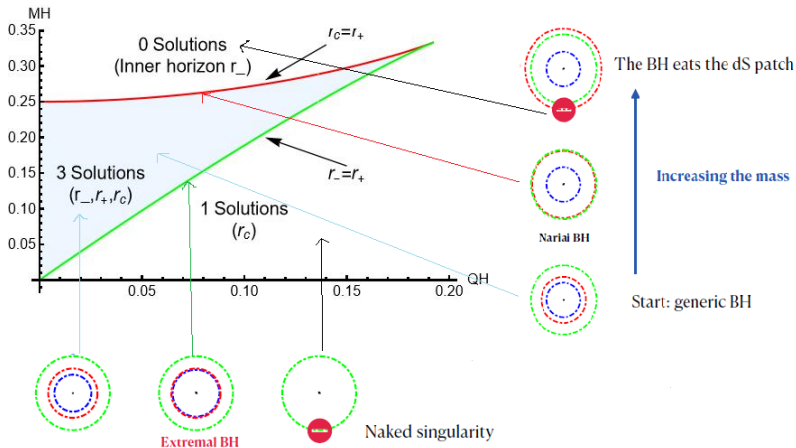
Increasing the charge:



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- ▶ In flat 5 dimensions $\Rightarrow r_- = r_+ \Rightarrow M^2 < \frac{3}{4} Q^2$.
- ▶ Now in De sitter space, by expansion of equation of green line

$$M^2 < \frac{3}{4} Q^2 + \frac{\sqrt{3}}{2\pi} Q^3 H^2 - \frac{3}{4\pi^2} Q^4 H^4 + \mathcal{O}(H^6)$$

Conclusion

- ▶ we investigated extremal black hole solutions in Einstein-Maxwell-Dilaton theory in extra dimensions.

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$$\alpha > \alpha_c \Rightarrow Q^2 e^{2\alpha\phi_0} = (1 + \alpha^2)M^2,$$

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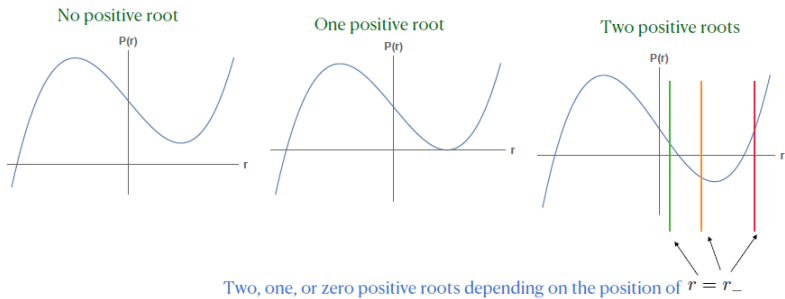
$$0 < \alpha < \alpha_c \Rightarrow \text{for } M^2 = \frac{3Q^2(2 - \alpha^2)}{8},$$

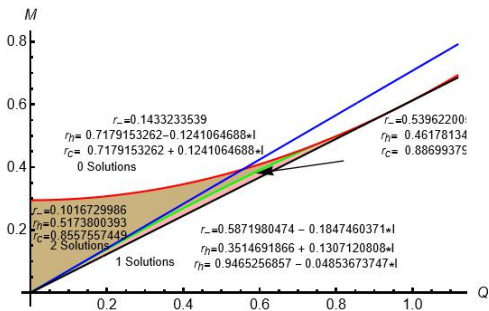
the metric is complex and the extremality is never reached

$$\alpha = 0 \Rightarrow M^2 < \frac{3}{4}Q^2 + \frac{\sqrt{3}}{2\pi}Q^3H^2 - \frac{3}{4\pi^2}Q^4H^4 + \mathcal{O}(H^6),$$

Thank you for your attention!

- ▶ In order to find the BH horizons, we look for the roots of $f(r)$,
- ★ We are interested only in solutions of $f(r)$ in the region $r > r_-$ outside the singularity





- The relation between the physical quantities mass M and Q^2 and the integration constants r_+ , r_- :

$$r_- = \frac{Q^2 (\alpha^2 + n - 3) \left(\frac{\sqrt{8(n-2)Q^2(\alpha^2 - n + 3) + M^2(n-3)^2} + M(n-3)}{(n-3)(n-2)} \right)^{3-n}}{(n-3)^2(n-2)}$$

$$r_+ = \frac{\left(\sqrt{M^2(n-3)^2 - 8(n-2)Q^2(-\alpha^2 + n - 3)} + M(n-3) \right)}{(n-3)(n-2)}$$

$$M^2 \geq \frac{8(n-2)Q^2(-\alpha^2 + n - 3)}{(n-3)^2}$$

