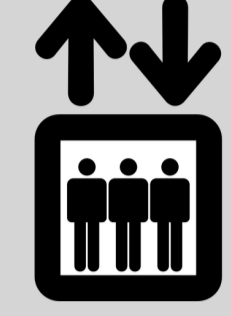


# The physics of the fragmentation region in heavy-ion collisions

Isobel Kolbé, Mawande Lushozi

Institute for Nuclear Theory, University of Washington, Seattle, Washington, USA  
 ikolbe@uw.edu



## Motivation

- The fragmentation region is interesting:
  - Access to densities  $\gtrsim 2$  \* nuclear saturation density ~ neutron star core density.
  - Give insight into cosmic-ray physics.
  - Not well studied - new frontiers!
  - The authors have previously studied [4] matter in the fragmentation region, but were forced to use a flat momentum spectrum for the gluon radiation. A more rigorous spectrum is needed i.o.t. fully model the behavior of fluids in this region.
- Technically interesting challenge:
  - The fragmentation region is dominated by soft (small- $x$ ) gluons, so we are dealing with the physics of the color-glass condensate (CGC).
  - But relativistic collisions necessarily also involve perturbative QCD physics.
  - How does one bridge the gap?

## The Basic Problem

- We will motivate an ansatz for gluon bremsstrahlung in the fragmentation region that properly connects the non-perturbative region (hereinafter the “classical result” [1] with the perturbative QCD region (hereinafter the “QCD” result) [2].
- The problem with the classical result:** The high- $k_T$  (momentum of the radiated gluon) limit of the classical result does not have the same fall-off as the QCD result (which must be correct at high- $k_T$ ).
- The problem with the QCD result:** It does not take into account the spectrum of transferred momentum  $q_\perp$ , which depends on the non-perturbative physics of the gluons in the nucleus (as an internal momentum,  $q_\perp$  is integrated out, but there should be a distribution of  $q_\perp$  in the non-perturbative limit, resulting from the multiple gluon interactions).
- We perform a perturbative calculation of bremsstrahlung for a scalar field which carries *classical* color charge (hereinafter the “scalar result”) which offers a bridge between the fully classical calculation and the fully fermionic QCD calculation.
- The scalar result suggests a modification of the classical result which leads to the correct high- $k_T$  fall-off but still includes the non-perturbative physics of the CGC

- In order to fully understand the **fragmentation region**, we need to know the momentum spectrum of **gluon bremsstrahlung**.
- Hard** because of **competing effects**:
  - The non-perturbative physics of colliding ultra-relativistic nuclei (gluons dominate,  $\therefore \exists$  **multiple gluon interactions**)
  - Perturbative** physics of high-frequency (high-momentum) radiation.
- Already known: non-perturbative limit [1,2], as well as the perturbative limit [3].
- The non-perturbative limit (called the “classical result” here) does not give the correct high-momentum behaviour.
- We present a scalar version of the full pQCD result that also uses a classical color charge to motivate a modification of the classical result.
- Having **modified the classical result**, we present an ansatz for gluon bremsstrahlung in the fragmentation region.
- This ansatz may now be used to model the behavior of fluids in the fragmentation region.

$$|\mathcal{M}_{sc.}|^2 = \frac{32m^2 P^2 g^6}{q_\perp^4} C(R)C_2^2(R) d_R x^2 (1-x)^2 \left[ m^2 \left( \frac{1}{D_A} - \frac{1}{D_B} \right)^2 - \frac{q_\perp^2}{D_A D_B} \right]$$

### What to notice:

The denominators  $D_A$  and  $D_B$  each go like  $k_T^2$ .  
 The QCD result therefore falls off like  $k_T^{-4}$ .

But the **classical** result has terms that fall off like  $k_T^{-2}$ .

The **scalar result**, though perturbative, is free of the clutter caused by the fermions in the QCD result

$$|\mathcal{M}_{bremm.}^{cl.}|^2 = -2x^2 M^2 \left( \frac{1}{D_A} + \frac{1-x}{D_B} \right)^2 + \frac{x}{D_A} - \frac{x(1-x)}{D_B} + x^2(1-x) \frac{q_\perp^2}{D_A D_B} \quad [1,2]$$

[3]

$$\begin{aligned} & |\mathcal{M}_A^{QCD}|^2 + |\mathcal{M}_B^{QCD}|^2 + 2 |\mathcal{M}_A^{QCD} \mathcal{M}_B^{QCD,*}| \\ &= \frac{4K}{d_R} C_R^2 C(R) \left[ 4x^2(x-1)M^2 \left( \frac{1}{D_A} - \frac{1}{D_B} \right)^2 + \frac{x^2(x^2 - 2x + 2)}{D_A D_B} q_\perp^2 \right] \end{aligned}$$

## Using the **scalar** result to justify a **perturbative** modification of the **classical** result:

Start with a change of notation to match the original classical result in [1,2], given there as:

$$\left. \frac{dN}{dy d^2k_T} \right|_{Class.} = \frac{g^2 C_F}{8\pi^3} \int \frac{d^2h}{(2\pi)^2} \tilde{S}(\mathbf{k} - \mathbf{h}) \left[ \frac{h^i}{h_T^2 + 2(k^-)^2} - \frac{k^i - \xi p^i}{|\mathbf{k} - \xi \mathbf{p}'|^2 + \xi^2 m^2} \right]^2$$

In this notation, the scalar result has the form

$$\left. \frac{dN}{dy d^2k_T} \right|_{Sc.} = \frac{g^2 C_R}{16\pi^3} \frac{1}{1-\xi} \left[ \frac{k^i}{k_T^2 + 2(k^-)^2} - (1+\xi) \frac{k^i - \xi p^i}{|\mathbf{k} - \xi \mathbf{p}'|^2 + \xi^2 m^2} \right]^2$$

The  $\tilde{S}(\mathbf{k} - \mathbf{h})$  contains the non-perturbative physics of the CGC (more precisely, it is the two-point correlation of Wilson lines).

In the McLerran-Venugopalan model,  $\tilde{S}(\mathbf{k} - \mathbf{h})$  is Gaussian.

In the high-energy limit (that is, in the limit where we expect the classical result to resemble the scalar result),  $\tilde{S}(\mathbf{k} - \mathbf{h})$  is a **delta function**, the integral over which sets  $\mathbf{h} \rightarrow \mathbf{k}$ . In this limit then, the classical result gives precisely the scalar result, except for the overall factor of  $(1-\xi)^{-1}$  and the factor of  $(1+\xi)$  in front of the second term.

We therefore posit a modification of the classical result, yielding an **ansatz for gluon bremsstrahlung in the fragmentation region**, as follows

$$\left. \frac{dN}{dy d^2k_T} \right|_{Ansatz} = \frac{g^2 C_F}{8\pi^3} \int \frac{d^2h}{(2\pi)^2} \tilde{S}(\mathbf{k} - \mathbf{h}) \left[ \frac{1}{1-\xi} \left[ \frac{h^i}{h_T^2 + 2(k^-)^2} - (1+\xi) \frac{k^i - \xi p^i}{|\mathbf{k} - \xi \mathbf{p}'|^2 + \xi^2 m^2} \right] \right]^2$$

Our Ansatz may now be used to model the momentum spectrum of bremsstrahlung in the fragmentation region.

The Ansatz has the following properties:

- In the small- $x$  limit ( $\xi \rightarrow 1$ ), the ansatz reduces to the classical result
- In the high energy limit ( $\tilde{S}(\mathbf{k} - \mathbf{h}) \rightarrow \delta(\mathbf{h} - \mathbf{k})$ ), the ansatz reduces to the perturbative (scalar result).

[1] K. Kajantie, Larry D. McLerran, and Risto Paatelainen. Gluon Radiation from a classical point particle II: dense gluon fields. *Phys. Rev. D*, 101(5):054012, 2020.

[2] Keijo Kajantie, Larry D. McLerran, and Risto Paatelainen. Gluon Radiation from a Classical Point Particle. *Phys. Rev. D*, 100(5):054011, 2019.

[3] Mawande Lushozi, Larry D. McLerran, Michał Praszalowicz, and Gongming Yu. Gluon Bremsstrahlung in Relativistic Heavy Ion Collisions. *Phys. Rev. C*, 102(3):034908, 2020.

[4] Isobel Kolbé, Mawande Lushozi, Larry D. McLerran, and Gongming Yu. Distribution of Nuclear Matter and Radiation in the Fragmentation Region. *Phys. Rev. C*, 103(4):044908, 2021.