

## Statistical and Thermal Models for Heavy Ion Collisions and Astrophysics B. Nemutudi<sup>a</sup>, C. Mare<sup>b</sup>, F. Mukosi<sup>c</sup>, K. Mafune<sup>c</sup>, T. Khumalo<sup>d,e</sup>, A. Muronga<sup>f</sup> University of Limpopo University of Pretoria University of Venda **Jniversity of the Witwatersrand** Themba I ABS Nelson Mandela Universit INTRODUCTION RESULTS article density for protons $\frac{g}{2\pi^2(\hbar c)^3} \int_{mc^2}^{\infty} \frac{\sqrt{\epsilon^2 + c^4 m^2} \epsilon d\epsilon}{e^{(\epsilon-\mu)/T} + 1}.$ 0.40 A thermal model in physics can be defined as a model that uses statistical mechanics, thermodynamics and To 0.35 ext 0.30 Thermal models can be used to describe particle multiplicity in heavy ion collisions as a function of the ast 을 0.25 Use <u>.</u> 0.20 che 0.15 0.10 erial view of Large Hadron Collider with other experiments like ALICE 0 75 100 125 150 175 200 Large Hadron Collider emperature in Me for pa a/. Physics Letters B 518 (2001) 41 on of experimental particle ratios and thermal model calcu-MeV., $\mu_b = 46$ MeV at $\sqrt{s} = 130$ GeV. tio Model calculation Experimental result $0.65\pm0.07$ 0.59 $0.08\pm0.01$ 0.09ĕ 0.6 - $1.00\pm0.02$ 1.00 $0.88\pm0.05$ 1.000.224 $0.149 \pm 0.02$ Redlich and S. Wheaton, Physics Letters B615 (2005) 50-54. We de Meson Dominated of rela 0.24 Fermi-0.22 Artist rendition of a neutron star merger $\Lambda < \pi >$ Source: https://www.ligo.caltech.edu/news/ligo20200106 0.20 🖺 0.18 · Meson Dominated Freeze-out <u>م</u> 0.16 Bose-E $K^{+}/\pi^{+}$ ŦŦ \_\_\_\_\_ 0.14 - $K/\pi$ 0.12 $dN_{\epsilon}$ 100 10 20 30 40 50 Pythor Vs\_NN (GeV) Open Quark density Built i $\frac{g_q}{2\pi^2} \frac{3}{2} T^3 \zeta(3)$ → hydrodynamics Interp Image: Courtesy of Jean Cleymans Good MatPlo 200 300 400 Temperature in MeV Fermi-Dirac Bose-Einstein Particles that have half- Integer-spin particles Th Do not obey Pauli's integer spin ACKNOWLEDGEMENTS Obey Pauli exclusion principle exclusion principle a Quantum effects considered • He de • W de $\bar{n}_j = \frac{1}{1 + e^{(\epsilon_j - \mu)/T}}$ $\bar{n}_j = \frac{1}{e^{(\epsilon_j - \mu)/T} - 1}$ ■ Tł hosting the internship. $N = \frac{gVm^{3/2}}{2^{1/2}\pi^2\hbar^3} \int \frac{\sqrt{\varepsilon} \,\mathrm{d}\varepsilon}{e^{(\varepsilon-\mu)/T} \pm 1}$ $N = \frac{gVm^{3/2}}{2^{1/2}\pi^2\hbar^3} \int \frac{\sqrt{\varepsilon} \,\mathrm{d}\varepsilon}{e^{(\varepsilon-\mu)/T}\pm 1}$ ea Th $\Omega = -\frac{2}{3} \frac{gVm^{3/2}}{2^{1/2}\pi^2\hbar^3} \int_{0}^{\infty} \frac{\varepsilon^{3/2} d\varepsilon}{e^{(\varepsilon-\mu)/T}\pm 1} d\varepsilon.$ $\Omega = -\frac{2}{3} \frac{gVm^{3/2}}{2^{1/2}\pi^2\hbar^3} \int \frac{e^{3/2} de}{e^{(e-\mu)/T} \pm 1}.$ predict realistic systems. Given enough time one can extend our model to investigate internship. other effects such as interactions, finite volume effects as well as studying the dependence of particle yields and ratios

statistical methods to describe the behaviour of a system with a variation in state variables.

- temperature T and baryon chemical potential  $\mu_b$ .
- The freeze-out stage is at thermal equilibrium hence the application of statistical physics.





Maxwell-Boltzmann

e.g. ideal gas

ignored)

Classical systems

(Quantum effects

 $\bar{n}_j = \frac{1}{e^{(E-\mu)/T}}$ 

Non-interacting and

distinguishable particles

<b>OBJECTIVES</b>	Р
understand the properties of nuclear matter under creme conditions in heavy ion collisions and prophysics	$0.4 = \frac{N}{V} = \frac{1}{2}$
e thermal and statistical models to understand the emical potential and temperature dependence of:	- 2.0 Farticle de
particle density	0.1 -
particle multiplicity and ratios	0.0
articles produced in heavy-ion collisions.	
	P. Braun-Munzinger Table 1: Comparison lations at $T = 174$ M Ratio $\overline{p/p}$ $\overline{p}/\pi^{-}$ $\pi^{-}/\pi$
<b>METHODOLOGY &amp; COMPUTATION</b>	$\frac{K^{-}/K}{K^{-}/\pi}$
erived the expressions for the particle density ativistic particles. -Dirac statistics: $E^2 = p^2c^2 + m^2c^4$ .	J. C., H. Oeschler, K.
$dN_{\epsilon} = \frac{gV}{2\pi^2(\hbar c)^3} \frac{\sqrt{\epsilon^2 + c^4 m^2} \epsilon d\epsilon}{e^{(\epsilon - \mu)/T} + 1} \qquad \qquad \frac{N}{V} = \frac{g}{2\pi^2(\hbar c)^3} \int_{mc^2}^{\infty} \frac{\sqrt{\epsilon^2 + c^4 m^2} \epsilon d\epsilon}{e^{(\epsilon - \mu)/T} + 1}.$	
Einstein statistics:	
$f_{\epsilon} = \frac{gV}{2\pi^2(\hbar c)^3} \frac{\sqrt{\epsilon^2 + c^4 m^2} \epsilon d\epsilon}{e^{(\epsilon - \mu)/T} - 1} \qquad \qquad \frac{N}{V} = \frac{g}{2\pi^2(\hbar c)^3} \int_{mc^2}^{\infty} \frac{\sqrt{\epsilon^2 + c^4 m^2 \epsilon}}{e^{(\epsilon - \mu)/T} - 1} d\epsilon$	
n used to solve the equations above	0 1
source, non-proprietary software.	
n support for scientific computing eg. SciPy and Numpy.	$n_a = \frac{1}{2}$
preted language = no need to compile.	6-
for prototyping and quickly implementing code.	ti k densit
otlib for graphs and visualization	۳ 2 -
<u>CONCLUSIONS</u>	0
he main aim was to understand and use different statistical odels to model particle yields, particle number densities and ratios. Sowever, the statistical models were studied and applied to escribe the last stage of Heavy Ion Collision (HIC). The managed to model the behaviour of quark and gluon ensities and this is compatible with results in literature. The models derived from undergraduate statistical ermodynamics knowledge are fitting experimental results on the Relativistic Heavy-Ion Collider (RHIC). This can be asily extended to LHC experimental results.	• V F • F F • T c
ie mouers are versaure as they can be renned to better	

on beam energy, volume, equilibration factors, spin and isospin



UNIVERSITY



We would like to thank National Institute for Theoretical Physics (NITheP) for funding the internship.

Prof. Azwinndini Muronga for the mentorship and

The Statistical and Thermal models for heavy ion collisions and astrophysics topical group members.

Mrs Rene Kotze for logistical arrangements during the