

# Structural properties of some defects in tin-dioxide (SnO<sub>2</sub>)

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**Abstract.** Tin-dioxide ceramics have been intensively studied in recent years because of the potential in sensing and fuel cells. The present work uses classical molecular dynamics simulations focused on the role of defects in tin-dioxide. The total energy of the NPT Hoover ensemble at various temperatures has been calculated in order to determine the effects of oxygen vacancy and Ti substitutional defect in tin-dioxide. The results obtained showed an energy increase with temperature which was constantly compared with experiments. The radial distribution functions for the structures suggest the transformation of anatase to rutile tin-dioxide at high temperature.

## 1. Introduction

Structural and thermodynamic properties of metal-oxide materials like tin oxide (SnO<sub>2</sub>) play crucial role as semiconducting sensors and electronic devices. The manifestation of intrinsic bulk and surface defect levels and the origin of extrinsic defect levels are of practical significance in the functioning property of these materials. The well-known reflectivity and conductivity property of SnO<sub>2</sub> are influenced by intrinsic native defects. In general, SnO<sub>2</sub> defects show up as oxygen vacancies. These introduce energy levels into the energy gap of the material causing the n-type conduction of states filled near the conduction band minimum. For effective use as sensors and semiconductors, these materials must have good transmittance of the visible light and resistivity well below 10<sup>-3</sup> Ωcm.

A great deal of experimental investigation has been performed on various forms of SnO<sub>2</sub> including the current nanoscale studies. Jones *et al* performed some scanning tunnelling microscopy measurements on the surface structure whilst Cox *et al* performed ion-scattering spectroscopy together with ultraviolet photoelectron measurements on the surface of SnO<sub>2</sub> to understand the oxygen vacancy defect in the conductivity of SnO<sub>2</sub>. Now advanced computer softwares based on empirical and first principle calculations make it possible to study the bulk, surface, and nanostructures of this material. Such development permits formation of an atomistic image of semiconduction and sensing mechanisms. Noteworthy, the first principle calculations are limited to zero temperature, although high temperature environments are responsible for a number of kinetic factors observed in thin films, surfaces, and nanoparticles. Subsequently, metal oxide sensors work at elevated temperatures, which also influence the impurity adsorption and catalytic processes in the material matrix. Through the molecular dynamics approach for the defects in SnO<sub>2</sub>, simulations at high temperatures can be achieved with ease.

The aim of this paper is to highlight some notable structural and thermodynamic properties obtained when anatase SnO<sub>2</sub> is subjected to oxygen vacancy defect and titanium substitutional defect. The empirical molecular dynamics using the Buckingham potentials has been used through the

simulations. Hopefully these results will provide an atomistic understanding of the role of defects in semiconducting sensors.

## 2. Computational details

Empirical molecular dynamics (MD) has been used to model the anatase SnO<sub>2</sub> and the related defects. A tetragonal space group I4<sub>1</sub>/amd SnO<sub>2</sub> unit cell has been repeated periodically in 3 dimensions to make a supercell with 192 Sn atoms and 384 O atoms. The oxygen parameter for anatase is 0.2066 Å and there are four Sn-O distance of 1.937 Å and two of 1.964 Å [7]. In the case of an oxygen vacancy defect on oxygen atom has been removed to make 383 O atoms. A Sn atom has also been replaced by a Ti atom in order for a substitutional Ti defect to materialize.

The material was modelled using the Buckingham potentials [8]. The DL-POLY package [8] has been used to perform all the empirical bond-order molecular dynamics calculations of SnO<sub>2</sub>. A supercell with a 7.29 cutoff, 576 atoms, and a sufficiently larger number grid points for the fast Fourier transformations (kmax1=6, kmax2=6, and kmax3=12) has been used for anatase SnO<sub>2</sub> throughout the calculations. The ewald convergence parameter of 0.3975 for anatase, on a Noose-Hoover NPT ensemble allowing the simulation supercell to change has been applied. The thermostat relaxation was set at 0.1 whilst the barostat was at 0.5. The simulation was allowed to run for more than 100 000 steps, with a simulation time step of 0.001ps. The controlled experimental crystal structure for anatase SnO<sub>2</sub> is according to Cromer and Herrington [7]. The material as is used for the molecular dynamics (MD) modeling is described by its lattice parameters as listed in Table 1, and a set of parameters required for the Buckingham potential are taken from AV Bandura *et al.* [9] and P Armstrong *et al.* [10]. Energy parameters were determined so as to reproduce the observed crystallographic structures of anatase SnO<sub>2</sub> and the accompanying dopants.

**Table 1. The lattice parameters and relative sites for O atom and Sn-O bonds [7].**

	this work	experiments
$a$ (Å)	3.8263	3.7845
$b$ (Å)	11.3194	9.5143
$u$	0.208	0.2066
Sn-O (Å)	4 x 2.1, 2 x 2.0	4 x 1.937, 2 x 1.964

## 3. Results and discussion

In Table I, the simulated lattice parameters of anatase SnO<sub>2</sub> are shown. The results show that the two-body Buckingham potential reproduce the crystallographic structures within the experimental results.

In this paper the thermodynamic properties of anatase SnO<sub>2</sub> together with oxygen vacancy defect (V<sub>O</sub>) in anatase SnO<sub>2</sub> and a Ti substitutional defect Ti<sub>Sn</sub> in anatase SnO<sub>2</sub> are being explored. The calculations have been performed above SnO<sub>2</sub> Debye temperature of 570 K by Tuerkes *et al.* [11] and 620 K by Bachmann *et al.* [12]. Subsequently the MD treats the motion of atoms classically; above the Debye temperature quantum mechanical effects can be neglected. So the average of the two, Debye temperatures was estimated at 595 K. Figure 1 shows the volume of anatase SnO<sub>2</sub>, V<sub>O</sub>, and Ti<sub>Sn</sub> as functions of temperature above the SnO<sub>2</sub> Debye temperature. From the plots it can be seen that Ti<sub>Sn</sub> curve has the lowest volume throughout the temperatures. This suggests that the presence of Ti substitutional defect in SnO<sub>2</sub> could assist in reduction of grain growth as argued by Rumyantseva *et al.* [13]. The volume thermal expansion coefficient for anatase SnO<sub>2</sub> was calculated to be 8.08 x 10<sup>-6</sup> K<sup>-1</sup>, of which differs by a few orders from the measured value of 11.7 x 10<sup>-6</sup> K<sup>-1</sup> by Percy and Morosin [14]. Of course it should also be noted that Percy and Morosin Raman measurements were done on a rutile SnO<sub>2</sub> structure.

The specific heat of anatase SnO<sub>2</sub> has also been calculated with the increasing volume. This is obtained from the temperature derivative of the total energy of the system. Figure 2 shows the said plot of energy against temperature for anatase SnO<sub>2</sub>, oxygen vacancy V<sub>O</sub>, substitutional Ti<sub>Sn</sub>. The

specific heat calculated is  $3.41k_B$  which differ by about 12% from the Dulong-Petit's law of solids at high temperatures. But even the measured values seem to be more offline with the value  $6.32k_B$  [15]. It can be seen on the plots that substitutional  $Ti_{Sn}$  has the lowest energy, which suggest a good dopant for enhanced conductivity. Likewise, the oxygen vacancy  $V_O$  and the anatase  $SnO_2$  curve intersect around 3000 K, which suggest an anticipated high temperature transition to the rutile  $SnO_2$  phase in agreement with the Fan and Reid [16] measurements.

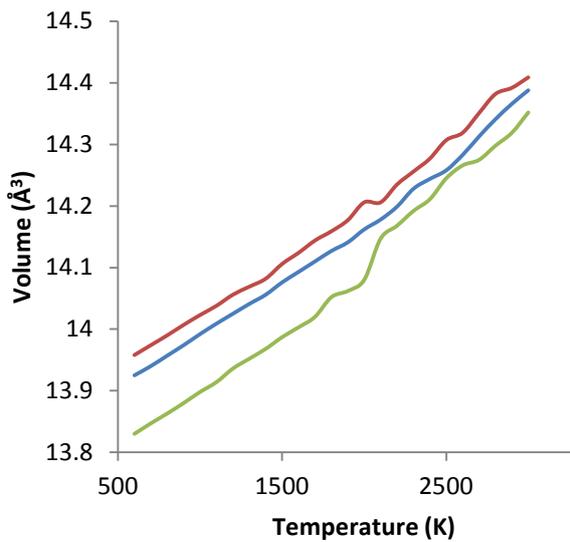


Figure 1. Volume as function of temperature for anatase  $SnO_2$ . blue curve is anatase  $SnO_2$ , red curve is  $V_O$ , and green curve is  $Ti_{Sn}$ .

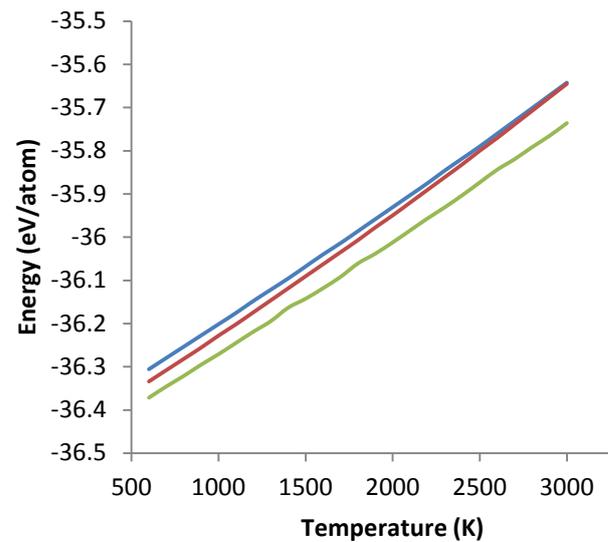


Figure 2. Energy as function of temperature for anatase  $SnO_2$ . blue curve is anatase  $SnO_2$ , red curve is  $V_O$ , and green curve is  $Ti_{Sn}$ .

In order to further check the stability of anatase  $SnO_2$  and its two defects at various temperatures within the Buckingham potential, the radial distribution functions of  $SnO_2$  at 300 K and 5000 K are presented in Figure 3 and 4. The peaks show the most probable distances between the atoms concerned. Furthermore, Figure 4 shows that at 5000 K, the peak positions for anatase  $SnO_2$  look the same as that of the rutile  $SnO_2$  at 300 K to 1000 K, with notable peak shifts in the rutile  $SnO_2$  [17]. The same form of peak positions is experienced with the substitutional  $Ti_{Sn}$ . This suggest a probable phase transition from anatase to rutile  $SnO_2$  in agreement with the first principle calculations of Yanlu Li *et al.* [18] and the experiments of Bachmann *et al.* [12] and Fan and Reid [16]. On the other hand these calculations show stability and transferability of the Buckingham potentials in structural analysis of materials.

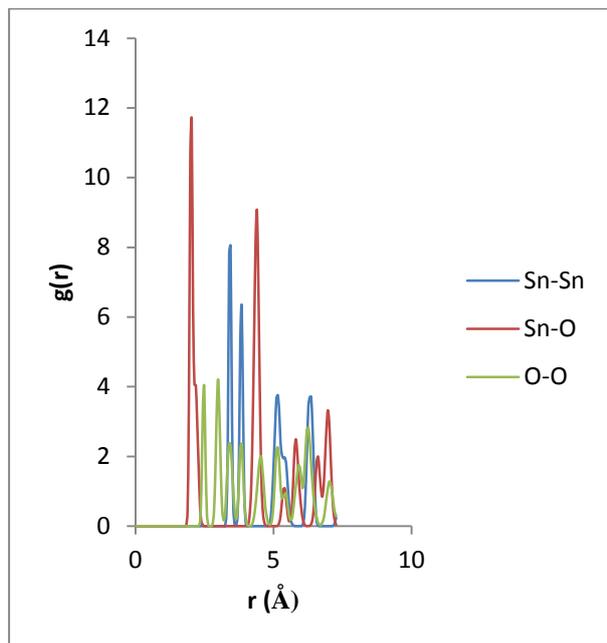


Figure 3. Radial distribution function of anatase SnO<sub>2</sub> at 300 K.

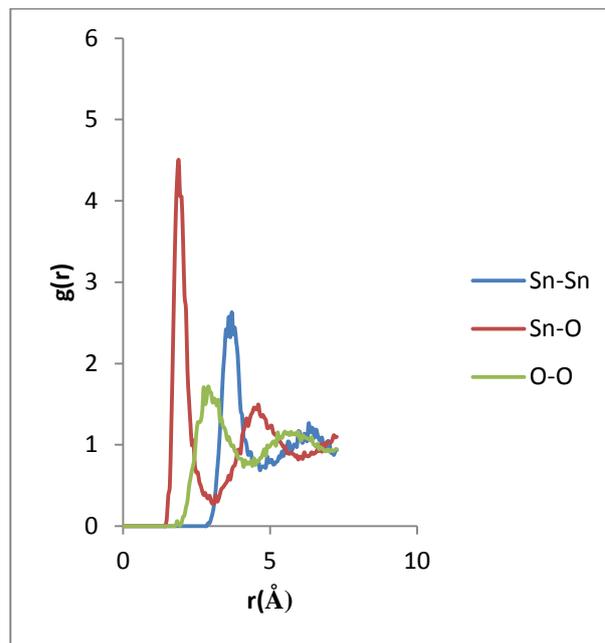


Figure 4. Radial distribution function of anatase SnO<sub>2</sub> at 5000 K.

#### 4. Conclusion

Anatase SnO<sub>2</sub>, oxygen vacancy in SnO<sub>2</sub> and Ti substitutional defect in SnO<sub>2</sub> have been modelled with molecular dynamics simulation using the Buckingham potential to investigate their structural and thermodynamic properties. The radial distribution curves and the energy-temperature graph suggest phase transformation at temperatures around 3000 K. Ti substitutional defect in anatase SnO<sub>2</sub> has the lowest energy and lowest volume which could improve its semiconducting properties at a controlled nanocrystalline growth. The volume thermal expansion coefficient is of the same order as the measured results. The specific heat capacity is of the same order with the Dulong-Petit law of solids at high temperatures.

#### 5. Acknowledgements

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