Modelling the Stellar Soft-photon Energy Density of Globular Clusters

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Presentation Outline

- I. Globular Clusters (GCs):
 - Millisecond pulsar (MSP) hosts
 - Recent gamma-ray observations
- 2. Inverse Compton (IC) scattering
- 3. Energy density profiles
 - Application to Terzan 5
- 4. Resulting IC-spectra
- 5. Model accuracy and improvements



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- → Spun-down pulsars gain angular momentum through mass-accretion
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✤ Fermi LAT and H.E.S.S. revealed GCs as sources of HE (>100 MeV) and VHE (>100 GeV) gamma-radiation

- \rightarrow for example, Terzan 5 (Ter5)
- \rightarrow 34 MSPs





(Freirre et al. 2011: *Fermi*-LAT gamma-ray (>100MeV) count map of NGC6642)



Inverse Compton (IC) scattering

Particles ejected by the MSP are accelerated to relativistic speeds (either in magnetosphere of MSP or due to relativistic shocks where pulsar winds collide).

Particles diffuse out of the globular cluster and interact with soft photons (CMB, IR, starlight).

***** The soft-photons are up-scattered as γ -rays in the TeV-band.





Inverse Compton (IC) scattering

* To calculate the IC-spectrum, consider the emissivity, given by Zhang *et al.* (2008):

 $Q_{Comp,j}(E_{\gamma},t) = 4\pi \int_{0}^{\infty} n_{j}(\epsilon,r)d\epsilon \int_{E_{e,thresh}}^{E_{e,max}} J_{e}(E_{e},t)F(\epsilon,E_{\gamma},E_{e})dE_{e}$

The component of interest for our purposes is

$$n_j(\epsilon) = \frac{15U_j}{\left(\pi kT_j\right)^4} \frac{\epsilon^2}{\left[e^{\left(\epsilon/kT_j\right)} - 1\right]}$$

 \clubsuit Energy density U_i

- Prominent stellar component in GCs
- Must decrease with increasing distance from cluster centre
- Our objective is to derive an energy density profile for the stellar/starlight component, and solve it for the case of Ter5.



✤First, we consider the contribution of a single star:

- Assume all stars in GCs radiate like blackbodies.
- Write down the result for the energy density contribution of a single star.
- Scale this result

 \odot down to compensate for the distance 'd' from the observer to the star,

 and up to account for the total radiating surface.

 $I_{\nu} = \frac{2h}{c^2} \frac{\nu^3}{e^{\frac{h\nu}{kT}} - 1}$ $u = \frac{4\pi}{c}I$ $u_s = \left(\frac{4\pi}{c}\right) \left(\frac{4\pi R^2}{4\pi d^2}\right) I$ $=\left(\frac{4\pi}{c}\right)\left(\frac{R^2}{d^2}\right)I$





We expand our result to include the contributions of all the stars:

• We approximate all the stars to have solar properties, and assume spherical symmetry.



 $u_s = \left(\frac{4\pi}{c}\right) \left(\frac{R^2}{d^2}\right) I$ $u(r) = \int u_s N(r') dr'$ $=\int u_s n(r')dV$ $=\int u_s \frac{\rho(r')}{\overline{m}} dV$

 $u(r) = \frac{4\pi R^2}{c} \frac{I}{\overline{m}} \int \frac{\rho(r')}{d^2} dV$



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$$u(r) = \frac{4\pi R^2}{c} \frac{I}{\overline{m}} \int \frac{\rho(r')}{d^2} dV$$

$$d^{2} = |\mathbf{r} - \mathbf{r}'|^{2}$$
$$= (r')^{2} + r^{2} - 2r'r\cos\theta$$

 $dV = (r')^2 \sin \theta \, dr' d\theta d\phi$

$$u(r) = \frac{8\pi^2 R^2}{c} \frac{I}{\overline{m}} \frac{1}{r} \int_0^{r_t} \rho(r') r' \ln \frac{|r'+r|}{|r'-r|} dr'$$





We consequently normalise the mass-density profile:

(Kuranov & Postnov 2006):

$$\rho(r) = \rho_0 \begin{cases} 1 & 0 < r < r_c \\ \left(\frac{r_c}{r}\right)^2 & r_c < r < r_h \\ \frac{(r_c r_h)^2}{r^4} & r_h < r < r_t \end{cases}$$



$$\begin{split} M_{total} &= N_{tot} \overline{m} \\ &= \int \rho(r) dV = 4\pi \int \rho(r) r^2 dr \end{split}$$





Comparison of energy densities for Ter5





-2 -1 Energy (TeV)

0

1

Scaled up with ~x3

-14

-4

-3







Concluding remarks





Predicted IC-spectrum:

- Provides a good fit to the H.E.S.S. data if scaled up by a factor 3
- N_star, N_MSP, eta and <E_dot> scaled up by ~1.3
- shows improvement
- ✤The error margins on u(r):
 - Propagated to the IC-spectrum in a linear fashion
 - H.E.S.S. data included within these error margins.

Improvements on the energy density profile:

- HR diagrams of GCs: Upper-limit masses, correct stellar relations.
- Surface brightness profiles
- Improvements on the IC-calculation:
 - Construct a cluster magnetic field profile
 - Use refined transport equations:
 - \odot Greater number of zones in radiation code without loss of stability

