

# GR Primer



Kevin Goldstein

School of Physics, University Witwatersrand

HDM2014



# Outline

- 1 Introduction
- 2 Geometry
  - Tensors
  - Curvature
- 3 General Relativity



# General Relativity

## GR Summary - 3 statements:

- Space-time is a curved pseudo-Riemannian manifold
- Matter curves space-time according to Einstein's equation:

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 8\pi GT_{\mu\nu}$$

- Free particles follow geodesics in space-time.
- We need to a bit of mathematics to appreciate the summary
- Work in units where  $c = 1$



# General Relativity

## GR Summary - 3 statements:

- Space-time is a curved pseudo-Riemannian manifold
- Matter curves space-time according to Einstein's equation:

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 8\pi GT_{\mu\nu}$$

- Free particles follow geodesics in space-time.
- We need to a bit of mathematics to appreciate the summary
- Work in units where  $c = 1$



# General Relativity

## GR Summary - 3 statements:

- Space-time is a curved pseudo-Riemannian manifold
- Matter curves space-time according to Einstein's equation:

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 8\pi GT_{\mu\nu}$$

- Free particles follow geodesics in space-time.
- We need to a bit of mathematics to appreciate the summary
- Work in units where  $c = 1$



# General Relativity

## GR Summary - 3 statements:

- Space-time is a curved pseudo-Riemannian manifold
- Matter curves space-time according to Einstein's equation:

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 8\pi GT_{\mu\nu}$$

- Free particles follow geodesics in space-time.
- We need to a bit of mathematics to appreciate the summary
- Work in units where  $c = 1$



# General Relativity

## GR Summary - 3 statements:

- Space-time is a curved pseudo-Riemannian manifold
- Matter curves space-time according to Einstein's equation:

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 8\pi GT_{\mu\nu}$$

- Free particles follow geodesics in space-time.
- We need to a bit of mathematics to appreciate the summary
- Work in units where  $c = 1$



# General Relativity

## GR Summary - 3 statements:

- Space-time is a curved pseudo-Riemannian manifold
- Matter curves space-time according to Einstein's equation:

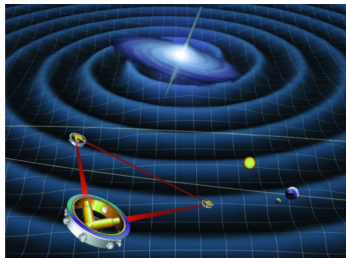
$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 8\pi GT_{\mu\nu}$$

- Free particles follow geodesics in space-time.
- We need to a bit of mathematics to appreciate the summary
- Work in units where  $c = 1$





## Recommended book



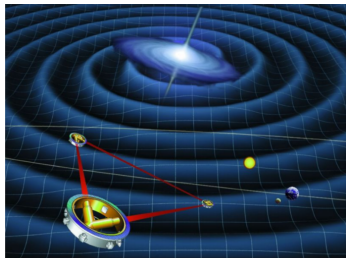
Credit: NASA

*A general relativity workbook* by Thomas Moore  
University Science Books, 2013

- A workbook you can work through yourself with effort and motivation



## Recommended book



Credit: NASA

*A general relativity workbook* by Thomas Moore  
University Science Books, 2013

- A workbook you can work through yourself with effort and motivation



# Riemannian Geometry

- Euclid's 5<sup>th</sup> postulate (E5)
  - parallel lines never meet
- Does not apply to spaces with curvature.



# Riemannian Geometry

- Euclid's 5<sup>th</sup> postulate (E5)
  - parallel lines never meet
- Does not apply to spaces with curvature.



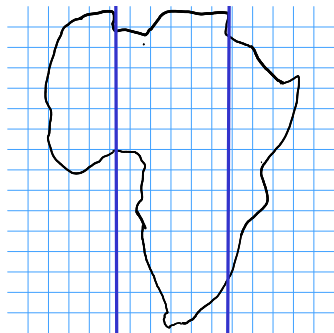
# Riemannian Geometry

- Euclid's 5<sup>th</sup> postulate (E5)
  - parallel lines never meet
- Does not apply to spaces with curvature.



# Riemannian Geometry

- Euclid's 5<sup>th</sup> postulate (E5)
  - parallel lines never meet
- Does not apply to spaces with curvature.

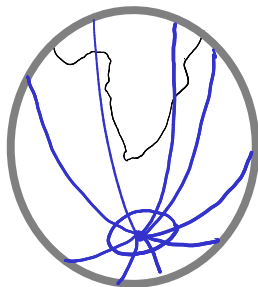


Parallel lines of longitude



# Riemannian Geometry

- Euclid's 5<sup>th</sup> postulate (E5)
  - parallel lines never meet
- Does not apply to spaces with curvature.



Parallel lines of longitude meet at the south pole

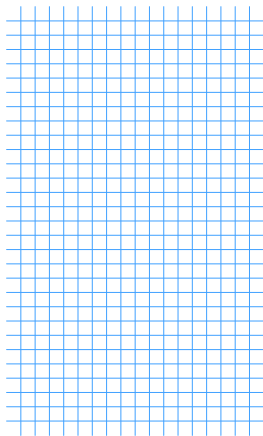


# Riemannian Geometry

- Euclid's 5<sup>th</sup> postulate (E5)
  - parallel lines never meetdoes not apply to spaces with curvature.
- The failure of E5 can be encoded in the **metric**
- The metric,  $ds^2$ , tells us the infinitesimal distance between points.
- Flat space metric:

$$ds^2 = dx^2 + dy^2 + dz^2$$

is equivalent to E5



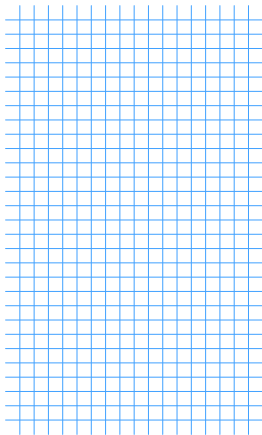


# Riemannian Geometry

- Euclid's 5<sup>th</sup> postulate (E5)
  - parallel lines never meetdoes not apply to spaces with curvature.
- The failure of E5 can be encoded in the **metric**
- The metric,  $ds^2$ , tells us the infinitesimal distance between points.
- Flat space metric:

$$ds^2 = dx^2 + dy^2 + dz^2$$

is equivalent to E5

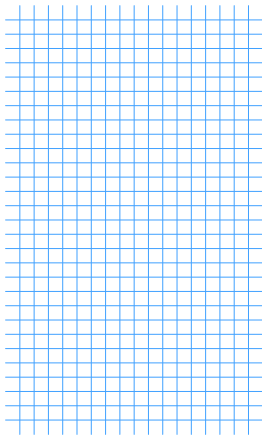


# Riemannian Geometry

- Euclid's 5<sup>th</sup> postulate (E5)
  - parallel lines never meetdoes not apply to spaces with curvature.
- The failure of E5 can be encoded in the **metric**
- The metric,  $ds^2$ , tells us the infinitesimal distance between points.
- Flat space metric:

$$ds^2 = dx^2 + dy^2 + dz^2$$

is equivalent to E5

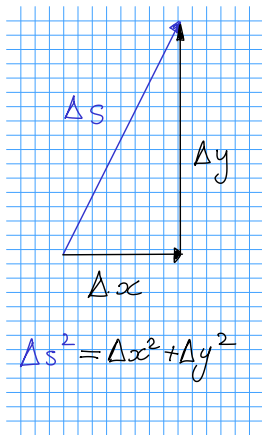


# Riemannian Geometry

- Euclid's 5<sup>th</sup> postulate (E5)
  - parallel lines never meetdoes not apply to spaces with curvature.
- The failure of E5 can be encoded in the **metric**
- The metric,  $ds^2$ , tells us the infinitesimal distance between points.
- Flat space metric:

$$ds^2 = dx^2 + dy^2 + dz^2$$

is equivalent to E5



# Metric

- Using the Einstein summation convention

$$ds^2 = \sum_{ij} g_{ij} dx^i dx^j \quad i, j \in 1, 2, 3$$

- In spherical coordinates

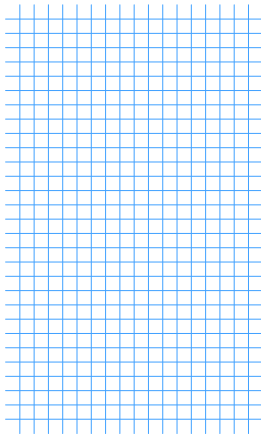
$$ds^2 = dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2$$

- $g_{ij} = \text{diag}(1, r^2, r^2 \sin^2 \theta)$

- In a general curved space the metric is a function of position,

$$g_{ij} = g_{ij}(\vec{x})$$

- Eg for the sphere,  
 $ds^2 = r^2(d\theta^2 + \sin^2 \theta d\phi^2)$



# Metric

- Using the Einstein summation convention

$$ds^2 = g_{ij} dx^i dx^j \quad i, j \in 1, 2, 3$$

- In spherical coordinates

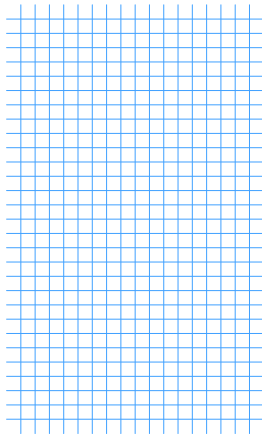
$$ds^2 = dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2$$

- $g_{ij} = \text{diag}(1, r^2, r^2 \sin^2 \theta)$

- In a general curved space the metric is a function of position,

$$g_{ij} = g_{ij}(\vec{x})$$

- Eg for the sphere,  
 $ds^2 = r^2(d\theta^2 + \sin^2 \theta d\phi^2)$



# Metric

- Using the Einstein summation convention

$$ds^2 = g_{ij} dx^i dx^j \quad i, j \in 1, 2, 3$$

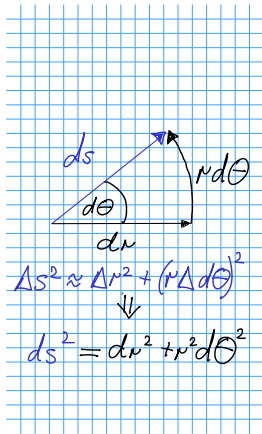
- In spherical coordinates

$$ds^2 = dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2$$

- $g_{ij} = \text{diag}(1, r^2, r^2 \sin^2 \theta)$

- In a general curved space the metric is a function of position,  $g_{ij} = g_{ij}(\vec{x})$

- Eg for the sphere,  
 $ds^2 = r^2(d\theta^2 + \sin^2 \theta d\phi^2)$



# Metric

- Using the Einstein summation convention

$$ds^2 = g_{ij} dx^i dx^j \quad i, j \in 1, 2, 3$$

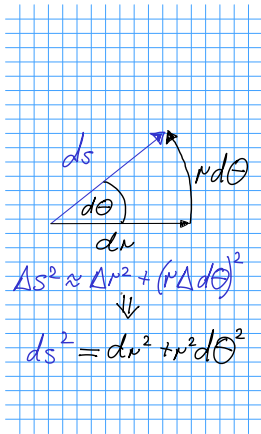
- In spherical coordinates

$$ds^2 = dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2$$

- $g_{ij} = \text{diag}(1, r^2, r^2 \sin^2 \theta)$

- In a general curved space the metric is a function of position,  $g_{ij} = g_{ij}(\vec{x})$

- Eg for the sphere,  
 $ds^2 = r^2(d\theta^2 + \sin^2 \theta d\phi^2)$



# Metric

- Using the Einstein summation convention

$$ds^2 = g_{ij} dx^i dx^j \quad i, j \in 1, 2, 3$$

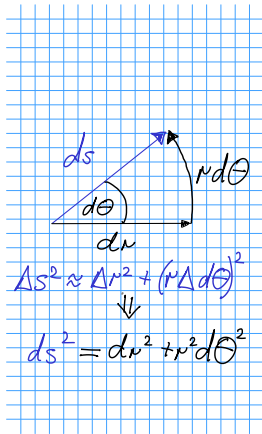
- In spherical coordinates

$$ds^2 = dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2$$

- $g_{ij} = \text{diag}(1, r^2, r^2 \sin^2 \theta)$

- In a general curved space the metric is a function of position,  $g_{ij} = g_{ij}(\vec{x})$

- Eg for the sphere,  
 $ds^2 = r^2(d\theta^2 + \sin^2 \theta d\phi^2)$





# Metric

- Using the Einstein summation convention

$$ds^2 = g_{ij} dx^i dx^j \quad i, j \in 1, 2, 3$$

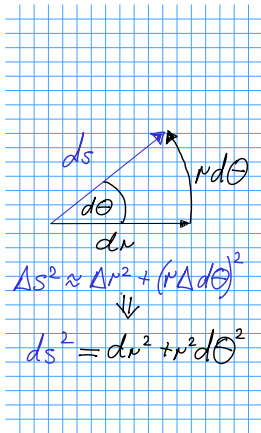
- In spherical coordinates

$$ds^2 = dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2$$

- $g_{ij} = \text{diag}(1, r^2, r^2 \sin^2 \theta)$

- In a general curved space the metric is a function of position,  $g_{ij} = g_{ij}(\vec{x})$

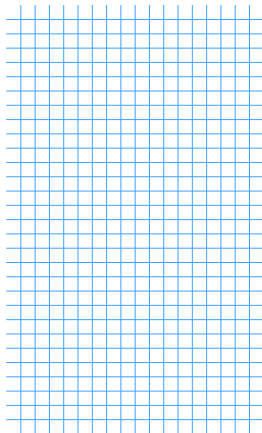
- Eg for the sphere,  
 $ds^2 = r^2(d\theta^2 + \sin^2 \theta d\phi^2)$



# Metric and Distance

- Metric = infinitesimal distance
- Consider a curve  $x^i(\lambda)$
- Integrate to get distance between points

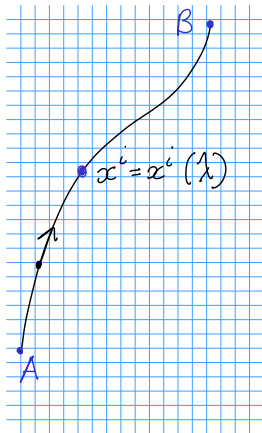
$$\begin{aligned} s_{AB} &= \int_A^B ds = \int_A^B d\lambda \frac{ds}{d\lambda} \\ &= \int_A^B d\lambda \sqrt{g_{ij} \frac{dx^i}{d\lambda} \frac{dx^j}{d\lambda}} \end{aligned}$$



# Metric and Distance

- Metric = infinitesimal distance
- Consider a curve  $x^i(\lambda)$
- Integrate to get distance between points

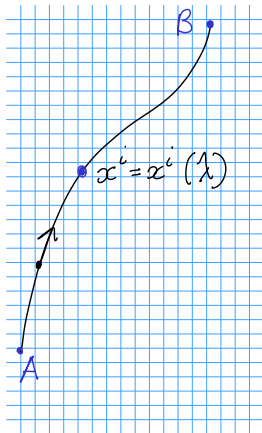
$$\begin{aligned} s_{AB} &= \int_A^B ds = \int_A^B d\lambda \frac{ds}{d\lambda} \\ &= \int_A^B d\lambda \sqrt{g_{ij} \frac{dx^i}{d\lambda} \frac{dx^j}{d\lambda}} \end{aligned}$$



# Metric and Distance

- Metric = infinitesimal distance
- Consider a curve  $x^i(\lambda)$
- Integrate to get distance between points

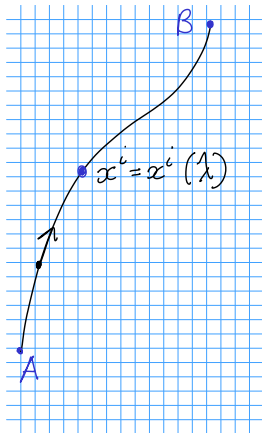
$$\begin{aligned} s_{AB} &= \int_A^B ds = \int_A^B d\lambda \frac{ds}{d\lambda} \\ &= \int_A^B d\lambda \sqrt{g_{ij} \frac{dx^i}{d\lambda} \frac{dx^j}{d\lambda}} \end{aligned}$$



# Metric and Distance

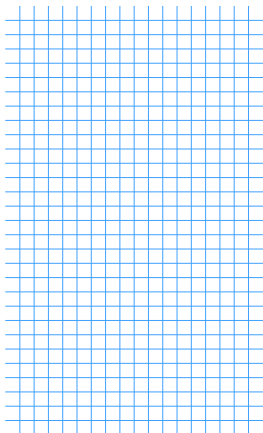
- Metric = infinitesimal distance
- Consider a curve  $x^i(\lambda)$
- Integrate to get distance between points

$$\begin{aligned} s_{AB} &= \int_A^B ds = \int_A^B d\lambda \frac{ds}{d\lambda} \\ &= \int_A^B d\lambda \sqrt{g_{ij} \frac{dx^i}{d\lambda} \frac{dx^j}{d\lambda}} \end{aligned}$$



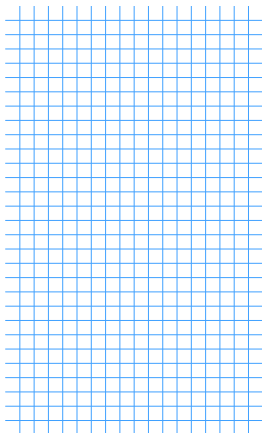
# Space-time

- speed of light constant  $\rightarrow$  special relativity
- $\rightarrow$  space-time interval same for all inertial observers:
  - $d\tau^2 = -dt^2 + dx^2 + dy^2 + dz^2 = g_{\mu\nu} dx^\nu dx^\mu \quad \mu, \nu \in 0, 1, 2, 3$
- The metric,  $g_{\mu\nu}$ , encodes the geometry of 4-d space-time
  - Flat space-time:  
 $g_{\mu\nu} = \text{diag}(-1, 1, 1, 1)$
- General curved pseudo-Riemannian geometry:  
 $g_{\mu\nu} = g_{\mu\nu}(x^\alpha)$



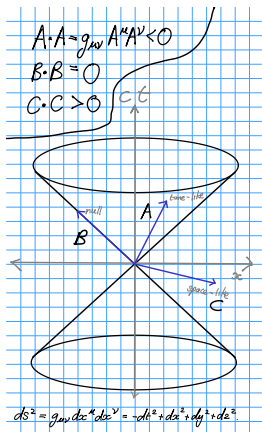
# Space-time

- speed of light constant  $\rightarrow$  special relativity
- $\rightarrow$  space-time interval same for all inertial observers:
  - $d\tau^2 = -dt^2 + dx^2 + dy^2 + dz^2 = g_{\mu\nu} dx^\nu dx^\mu \quad \mu, \nu \in 0, 1, 2, 3$
- The metric,  $g_{\mu\nu}$ , encodes the geometry of 4-d space-time
  - Flat space-time:  
 $g_{\mu\nu} = \text{diag}(-1, 1, 1, 1)$
- General curved pseudo-Riemannian geometry:  
 $g_{\mu\nu} = g_{\mu\nu}(x^\alpha)$



# Space-time

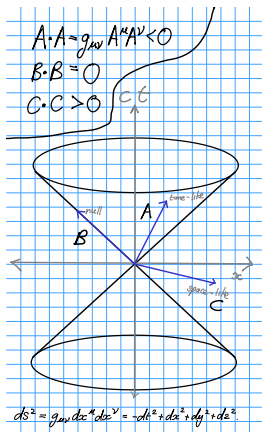
- speed of light constant  $\rightarrow$  special relativity
- $\rightarrow$  space-time interval same for all inertial observers:
  - $d\tau^2 = -dt^2 + dx^2 + dy^2 + dz^2 = g_{\mu\nu} dx^\nu dx^\mu \quad \mu, \nu \in 0, 1, 2, 3$
- The metric,  $g_{\mu\nu}$ , encodes the geometry of 4-d space-time
  - Flat space-time:  $g_{\mu\nu} = \text{diag}(-1, 1, 1, 1)$
- General curved pseudo-Riemannian geometry:  $g_{\mu\nu} = g_{\mu\nu}(x^\alpha)$





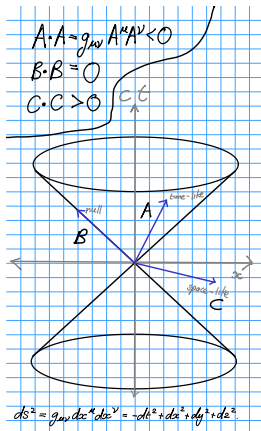
# Space-time

- speed of light constant  $\rightarrow$  special relativity
- $\rightarrow$  space-time interval same for all inertial observers:
  - $d\tau^2 = -dt^2 + dx^2 + dy^2 + dz^2 = g_{\mu\nu} dx^\nu dx^\mu \quad \mu, \nu \in 0, 1, 2, 3$
- The metric,  $g_{\mu\nu}$ , encodes the geometry of 4-d space-time
  - Flat space-time:  $g_{\mu\nu} = \text{diag}(-1, 1, 1, 1)$
- General curved pseudo-Riemannian geometry:  
 $g_{\mu\nu} = g_{\mu\nu}(x^\alpha)$



# Space-time

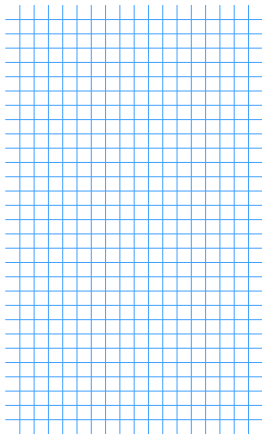
- speed of light constant  $\rightarrow$  special relativity
- $\rightarrow$  space-time interval same for all inertial observers:
  - $d\tau^2 = -dt^2 + dx^2 + dy^2 + dz^2 = g_{\mu\nu} dx^\nu dx^\mu \quad \mu, \nu \in 0, 1, 2, 3$
- The metric,  $g_{\mu\nu}$ , encodes the geometry of 4-d space-time
  - Flat space-time:  $g_{\mu\nu} = \text{diag}(-1, 1, 1, 1)$
- General curved pseudo-Riemannian geometry:  $g_{\mu\nu} = g_{\mu\nu}(x^\alpha)$



# Metric and Proper-time

- Metric in space-time also tells us how fast clocks run in space-time
- Consider a time-like curve  $x^\mu(\lambda)$
- Integrate to get proper time measured by an observer

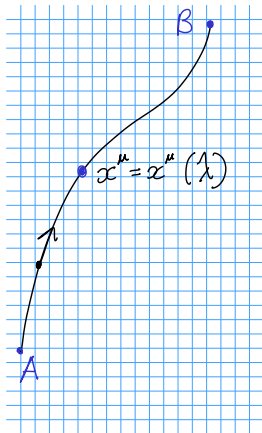
$$\begin{aligned}\tau_{AB} &= \int_A^B \sqrt{-ds^2} \\ &= \int_A^B d\lambda \sqrt{-g_{\mu\nu} \frac{dx^\mu}{d\lambda} \frac{dx^\nu}{d\lambda}}\end{aligned}$$



# Metric and Proper-time

- Metric in space-time also tells us how fast clocks run in space-time
- Consider a time-like curve  $x^\mu(\lambda)$
- Integrate to get proper time measured by an observer

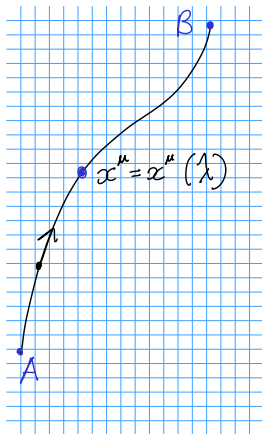
$$\begin{aligned}\tau_{AB} &= \int_A^B \sqrt{-ds^2} \\ &= \int_A^B d\lambda \sqrt{-g_{\mu\nu} \frac{dx^\mu}{d\lambda} \frac{dx^\nu}{d\lambda}}\end{aligned}$$



# Metric and Proper-time

- Metric in space-time also tells us how fast clocks run in space-time
- Consider a time-like curve  $x^\mu(\lambda)$
- Integrate to get proper time measured by an observer

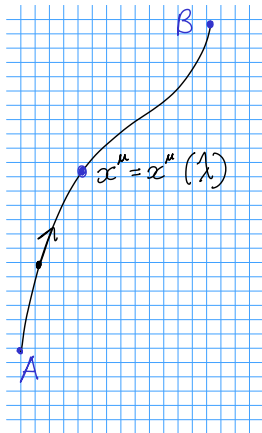
$$\begin{aligned}\tau_{AB} &= \int_A^B \sqrt{-ds^2} \\ &= \int_A^B d\lambda \sqrt{-g_{\mu\nu} \frac{dx^\mu}{d\lambda} \frac{dx^\nu}{d\lambda}}\end{aligned}$$



# Metric and Proper-time

- Metric in space-time also tells us how fast clocks run in space-time
- Consider a time-like curve  $x^\mu(\lambda)$
- Integrate to get proper time measured by an observer

$$\begin{aligned}\tau_{AB} &= \int_A^B \sqrt{-ds^2} \\ &= \int_A^B d\lambda \sqrt{-g_{\mu\nu} \frac{dx^\mu}{d\lambda} \frac{dx^\nu}{d\lambda}}\end{aligned}$$



# Metric Tensor

- Changing coordinates  $x^\mu \rightarrow x^{\bar{\alpha}}$ , distance between points must remain the same

⇒

$$g_{\bar{\alpha}\bar{\beta}} = \frac{\partial x^\mu}{\partial x^{\bar{\alpha}}} \frac{\partial x^\nu}{\partial x^{\bar{\beta}}} g_{\mu\nu}$$

- Coordinate transformation property of metric
- E.g. in flat space, Cartesian  $\rightarrow$  Spherical coordinates,

$$\text{diag}(1, 1, 1) \rightarrow \text{diag}(1, r^2, r^2 \sin^2 \theta)$$



# Metric Tensor

- Changing coordinates  $x^\mu \rightarrow x^{\bar{\alpha}}$ , distance between points must remain the same

$$\Rightarrow ds^2 = g_{\mu\nu} dx^\mu dx^\nu = g_{\bar{\alpha}\bar{\beta}} dx^{\bar{\alpha}} dx^{\bar{\beta}}$$

$\Rightarrow$

$$g_{\bar{\alpha}\bar{\beta}} = \frac{\partial x^\mu}{\partial x^{\bar{\alpha}}} \frac{\partial x^\nu}{\partial x^{\bar{\beta}}} g_{\mu\nu}$$

- Coordinate transformation property of metric
- E.g. in flat space, Cartesian  $\rightarrow$  Spherical coordinates,

$$\text{diag}(1, 1, 1) \rightarrow \text{diag}(1, r^2, r^2 \sin^2 \theta)$$





# Metric Tensor

- Changing coordinates  $x^\mu \rightarrow x^{\bar{\alpha}}$ , distance between points must remain the same

$$\Rightarrow ds^2 = g_{\mu\nu} dx^\mu dx^\nu = g_{\mu\nu} \left( \frac{\partial x^\mu}{\partial x^{\bar{\alpha}}} dx^{\bar{\alpha}} \right) \left( \frac{\partial x^\nu}{\partial x^{\bar{\beta}}} dx^{\bar{\beta}} \right) = g_{\bar{\alpha}\bar{\beta}} dx^{\bar{\alpha}} dx^{\bar{\beta}}$$

$\Rightarrow$

$$g_{\bar{\alpha}\bar{\beta}} = \frac{\partial x^\mu}{\partial x^{\bar{\alpha}}} \frac{\partial x^\nu}{\partial x^{\bar{\beta}}} g_{\mu\nu}$$

- Coordinate transformation property of metric
- E.g. in flat space, Cartesian  $\rightarrow$  Spherical coordinates,

$$\text{diag}(1, 1, 1) \rightarrow \text{diag}(1, r^2, r^2 \sin^2 \theta)$$



# Metric Tensor

- Changing coordinates  $x^\mu \rightarrow x^{\bar{\alpha}}$ , distance between points must remain the same

$$\Rightarrow ds^2 = g_{\mu\nu} dx^\mu dx^\nu = g_{\mu\nu} \left( \frac{\partial x^\mu}{\partial x^{\bar{\alpha}}} dx^{\bar{\alpha}} \right) \left( \frac{\partial x^\nu}{\partial x^{\bar{\beta}}} dx^{\bar{\beta}} \right) = g_{\bar{\alpha}\bar{\beta}} dx^{\bar{\alpha}} dx^{\bar{\beta}}$$

$\Rightarrow$

$$g_{\bar{\alpha}\bar{\beta}} = \frac{\partial x^\mu}{\partial x^{\bar{\alpha}}} \frac{\partial x^\nu}{\partial x^{\bar{\beta}}} g_{\mu\nu}$$

- Coordinate transformation property of metric
- E.g. in flat space, Cartesian  $\rightarrow$  Spherical coordinates,

$$\text{diag}(1, 1, 1) \rightarrow \text{diag}(1, r^2, r^2 \sin^2 \theta)$$



# Metric Tensor

- Changing coordinates  $x^\mu \rightarrow x^{\bar{\alpha}}$ , distance between points must remain the same

$$\Rightarrow ds^2 = g_{\mu\nu} dx^\mu dx^\nu = g_{\mu\nu} \left( \frac{\partial x^\mu}{\partial x^{\bar{\alpha}}} dx^{\bar{\alpha}} \right) \left( \frac{\partial x^\nu}{\partial x^{\bar{\beta}}} dx^{\bar{\beta}} \right) = g_{\bar{\alpha}\bar{\beta}} dx^{\bar{\alpha}} dx^{\bar{\beta}}$$

$\Rightarrow$

$$g_{\bar{\alpha}\bar{\beta}} = \frac{\partial x^\mu}{\partial x^{\bar{\alpha}}} \frac{\partial x^\nu}{\partial x^{\bar{\beta}}} g_{\mu\nu}$$

- Coordinate transformation property of metric
- E.g. in flat space, Cartesian  $\rightarrow$  Spherical coordinates,

$$\text{diag}(1, 1, 1) \rightarrow \text{diag}(1, r^2, r^2 \sin^2 \theta)$$



# Metric Tensor

- Changing coordinates  $x^\mu \rightarrow x^{\bar{\alpha}}$ , distance between points must remain the same

$$\Rightarrow ds^2 = g_{\mu\nu} dx^\mu dx^\nu = g_{\mu\nu} \left( \frac{\partial x^\mu}{\partial x^{\bar{\alpha}}} dx^{\bar{\alpha}} \right) \left( \frac{\partial x^\nu}{\partial x^{\bar{\beta}}} dx^{\bar{\beta}} \right) = g_{\bar{\alpha}\bar{\beta}} dx^{\bar{\alpha}} dx^{\bar{\beta}}$$

$\Rightarrow$

$$g_{\bar{\alpha}\bar{\beta}} = \frac{\partial x^\mu}{\partial x^{\bar{\alpha}}} \frac{\partial x^\nu}{\partial x^{\bar{\beta}}} g_{\mu\nu}$$

- Coordinate transformation property of metric
- E.g. in flat space, Cartesian  $\rightarrow$  Spherical coordinates,

$$\text{diag}(1, 1, 1) \rightarrow \text{diag}(1, r^2, r^2 \sin^2 \theta)$$



# General Tensors

General Tensors transform according to the rule

$$S^{\bar{\alpha}\dots}_{\bar{\beta}\bar{\gamma}\dots}(\bar{x}) = \left(\frac{\partial x^{\bar{\alpha}}}{\partial x^{\mu}}\right) \cdots \left(\frac{\partial x^{\nu}}{\partial x^{\bar{\beta}}}\right) \left(\frac{\partial x^{\rho}}{\partial x^{\bar{\gamma}}}\right) \cdots S^{\mu\dots}_{\nu\rho\dots}(x)$$

- To remember the rule:
  - “Conservation of indices ” : number of free (unsummed) indices must balance
  - dummy (summed) indices in pairs.
- Definition: An  $\binom{n}{m}$  tensor has n upper and m lower indices.
- A scalar is a  $\binom{0}{0}$  tensor.



# General Tensors

General Tensors transform according to the rule

$$S^{\bar{\alpha}\dots}_{\bar{\beta}\bar{\gamma}\dots}(\bar{x}) = \left(\frac{\partial x^{\bar{\alpha}}}{\partial x^{\mu}}\right) \cdots \left(\frac{\partial x^{\nu}}{\partial x^{\bar{\beta}}}\right) \left(\frac{\partial x^{\rho}}{\partial x^{\bar{\gamma}}}\right) \cdots S^{\mu\dots}_{\nu\rho\dots}(x)$$

- To remember the rule:
  - “Conservation of indices ” : number of free (unsummed) indices must balance
  - dummy (summed) indices in pairs.
- Definition: An  $\binom{n}{m}$  tensor has n upper and m lower indices.
- A scalar is a  $\binom{0}{0}$  tensor.



# General Tensors

General Tensors transform according to the rule

$$S^{\bar{\alpha}\dots}_{\bar{\beta}\bar{\gamma}\dots}(\bar{x}) = \left(\frac{\partial x^{\bar{\alpha}}}{\partial x^{\mu}}\right) \cdots \left(\frac{\partial x^{\nu}}{\partial x^{\bar{\beta}}}\right) \left(\frac{\partial x^{\rho}}{\partial x^{\bar{\gamma}}}\right) \cdots S^{\mu\dots}_{\nu\rho\dots}(x)$$

- To remember the rule:
  - “Conservation of indices ” : number of free (unsummed) indices must balance
  - dummy (summed) indices in pairs.
- Definition: An  $\binom{n}{m}$  tensor has n upper and m lower indices.
- A scalar is a  $\binom{0}{0}$  tensor.



# General Tensors

General Tensors transform according to the rule

$$S^{\bar{\alpha}\dots}_{\bar{\beta}\bar{\gamma}\dots}(\bar{x}) = \left(\frac{\partial x^{\bar{\alpha}}}{\partial x^{\mu}}\right) \dots \left(\frac{\partial x^{\nu}}{\partial x^{\bar{\beta}}}\right) \left(\frac{\partial x^{\rho}}{\partial x^{\bar{\gamma}}}\right) \dots S^{\mu\dots}_{\nu\rho\dots}(x)$$

- To remember the rule:
  - “Conservation of indices ” : number of free (unsummed) indices must balance
  - dummy (summed) indices in pairs.
- Definition: An  $\binom{n}{m}$  tensor has n upper and m lower indices.
- A scalar is a  $\binom{0}{0}$  tensor.





# General Tensors

General Tensors transform according to the rule

$$S^{\bar{\alpha}\dots}_{\bar{\beta}\bar{\gamma}\dots}(\bar{x}) = \left(\frac{\partial x^{\bar{\alpha}}}{\partial x^{\mu}}\right) \dots \left(\frac{\partial x^{\nu}}{\partial x^{\bar{\beta}}}\right) \left(\frac{\partial x^{\rho}}{\partial x^{\bar{\gamma}}}\right) \dots S^{\mu\dots}_{\nu\rho\dots}(x)$$

- To remember the rule:
  - “Conservation of indices ” : number of free (unsummed) indices must balance
  - dummy (summed) indices in pairs.
- Definition: An  $\binom{n}{m}$  tensor has n upper and m lower indices.
- A scalar is a  $\binom{0}{0}$  tensor.



# Why Tensors?

- Principle of general relativity: Laws of physics should look the same in *all* coordinate systems.
- Tensor equations look the same in all coordinates:

$$S^{\bar{\alpha}\dots}_{\bar{\beta}\bar{\gamma}\dots} = 0 \Rightarrow S^{\mu\dots}_{\nu\lambda} = 0$$

⇒ Tensors are the natural language of classical physics.



# Why Tensors?

- Principle of general relativity: Laws of physics should look the same in *all* coordinate systems.
- Tensor equations look the same in all coordinates:

$$S^{\bar{\alpha}\dots}_{\bar{\beta}\bar{\gamma}\dots} = 0 \Rightarrow S^{\mu\dots}_{\nu\lambda} = 0$$

⇒ Tensors are the natural language of classical physics.



# Why Tensors?

- Principle of general relativity: Laws of physics should look the same in *all* coordinate systems.
- Tensor equations look the same in all coordinates:

$$S^{\bar{\alpha}\dots}_{\bar{\beta}\bar{\gamma}\dots} = 0 \Rightarrow S^{\mu\dots}_{\nu\lambda} = 0$$

⇒ Tensors are the natural language of classical physics.



# Why Tensors?

- Principle of general relativity: Laws of physics should look the same in *all* coordinate systems.
- Tensor equations look the same in all coordinates:

$$S^{\bar{\alpha}\dots}_{\bar{\beta}\bar{\gamma}\dots} = 0 \Rightarrow S^{\mu\dots}_{\nu\lambda} = 0$$

⇒ Tensors are the natural language of classical physics.



# Tensor gymnastics

Changing tensor type using the metric

$$A_\mu = g_{\mu\nu} A^\nu$$

$\binom{1}{0}$ -tensor (vector)  $\rightarrow$   $\binom{0}{1}$ -tensor (co-vector)

Definition: Inverse metric  $g^{\mu\nu}$ :

Matrix inverse of the metric  $g^{\mu\nu} g_{\nu\alpha} = \delta_\alpha^\mu$



# Tensor gymnastics

## Lowering index

$$A_\mu = g_{\mu\nu} A^\nu$$

$\binom{1}{0}$ -tensor (vector)  $\rightarrow$   $\binom{0}{1}$ -tensor (co-vector)

Definition: Inverse metric  $g^{\mu\nu}$ :

Matrix inverse of the metric  $g^{\mu\nu} g_{\nu\alpha} = \delta_\alpha^\mu$



# Tensor gymnastics

## Lowering index

$$A^{\alpha\dots}_{\mu\beta\dots} = g_{\mu\nu} A^{\alpha\dots\nu}_{\beta\dots}$$

$\binom{n}{m}$ -tensor (vector)  $\rightarrow$   $\binom{n-1}{m+1}$ -tensor

Definition: Inverse metric  $g^{\mu\nu}$ :

Matrix inverse of the metric  $g^{\mu\nu} g_{\nu\alpha} = \delta^{\mu}_{\alpha}$





# Tensor gymnastics

## Lowering index

$$A^{\alpha\dots}_{\mu\beta\dots} = g_{\mu\nu} A^{\alpha\dots\nu}_{\beta\dots}$$

$\binom{n}{m}$ -tensor (vector)  $\rightarrow$   $\binom{n-1}{m+1}$ -tensor

Definition: Inverse metric  $g^{\mu\nu}$ :

Matrix inverse of the metric  $g^{\mu\nu} g_{\nu\alpha} = \delta_{\alpha}^{\mu}$



# Tensor gymnastics

## Lowering index

$$A^{\alpha\dots}_{\mu\beta\dots} = g_{\mu\nu} A^{\alpha\dots\nu}_{\beta\dots}$$

$\binom{n}{m}$ -tensor (vector)  $\rightarrow$   $\binom{n-1}{m+1}$ -tensor

## Definition: Inverse metric $g^{\mu\nu}$ :

Matrix inverse of the metric  $g^{\mu\nu} g_{\nu\alpha} = \delta^{\mu}_{\alpha}$

## Changing tensor type using the Inverse metric

$$A^{\mu} = g^{\mu\nu} A_{\nu}$$

$\binom{0}{1}$ -tensor (co-vector)  $\rightarrow$   $\binom{1}{0}$ -tensor (vector)



# Tensor gymnastics

## Lowering index

$$A^{\alpha\dots}_{\mu\beta\dots} = g_{\mu\nu} A^{\alpha\dots\nu}_{\beta\dots}$$

$\binom{n}{m}$ -tensor (vector)  $\rightarrow$   $\binom{n-1}{m+1}$ -tensor

## Definition: Inverse metric $g^{\mu\nu}$ :

Matrix inverse of the metric  $g^{\mu\nu} g_{\nu\alpha} = \delta^{\mu}_{\alpha}$

## Raising Index

$$A^{\mu} = g^{\mu\nu} A_{\nu}$$

$\binom{0}{1}$ -tensor (co-vector)  $\rightarrow$   $\binom{1}{0}$ -tensor (vector)



# Tensor gymnastics

## Lowering index

$$A^{\alpha\dots}_{\mu\beta\dots} = g_{\mu\nu} A^{\alpha\dots\nu}_{\beta\dots}$$

$\binom{n}{m}$ -tensor (vector)  $\rightarrow$   $\binom{n-1}{m+1}$ -tensor

## Definition: Inverse metric $g^{\mu\nu}$ :

Matrix inverse of the metric  $g^{\mu\nu} g_{\nu\alpha} = \delta^{\mu}_{\alpha}$

## Raising Index

$$A_{\alpha\dots}{}^{\mu\beta\dots} = g^{\mu\nu} A_{\alpha\dots\nu}{}^{\beta\dots}$$

$\binom{n}{m}$ -tensor (vector)  $\rightarrow$   $\binom{n+1}{m-1}$ -tensor



# Covariant derivative

## Notation

$$\frac{\partial}{\partial x^\mu} A = \partial_\mu A = A_{,\mu}$$

- Partial derivatives of scalar fields are  $\binom{0}{1}$ -tensors

$$\phi_{,\mu} = \left( \frac{\partial x^{\bar{\alpha}}}{\partial x^\mu} \right) \phi_{,\bar{\alpha}}$$

- Partial derivatives of other tensor fields are not tensors

$$A^\mu_{,\nu} = \frac{\partial x^{\bar{\alpha}}}{\partial x^\nu} \frac{\partial}{\partial x^{\bar{\alpha}}} \left( \frac{\partial x^\mu}{\partial x^{\bar{\beta}}} A^{\bar{\beta}} \right) = \frac{\partial x^{\bar{\alpha}}}{\partial x^\nu} \frac{\partial x^\mu}{\partial x^{\bar{\beta}}} A^{\bar{\beta}}_{,\bar{\alpha}} + \frac{\partial x^{\bar{\alpha}}}{\partial x^\nu} \frac{\partial^2 x^\mu}{\partial x^{\bar{\beta}} \partial x^{\bar{\alpha}}} A^{\bar{\beta}}$$



# Covariant derivative

## Notation

$$\frac{\partial}{\partial x^\mu} A = \partial_\mu A = A_{,\mu}$$

- Partial derivatives of scalar fields are  $\binom{0}{1}$ -tensors

$$\phi_{,\mu} = \left( \frac{\partial x^{\bar{\alpha}}}{\partial x^\mu} \right) \phi_{,\bar{\alpha}}$$

- Partial derivatives of other tensor fields are not tensors

$$A^\mu_{,\nu} = \frac{\partial x^{\bar{\alpha}}}{\partial x^\nu} \frac{\partial}{\partial x^{\bar{\alpha}}} \left( \frac{\partial x^\mu}{\partial x^{\bar{\beta}}} A^{\bar{\beta}} \right) = \frac{\partial x^{\bar{\alpha}}}{\partial x^\nu} \frac{\partial x^\mu}{\partial x^{\bar{\beta}}} A^{\bar{\beta}}_{,\bar{\alpha}} + \frac{\partial x^{\bar{\alpha}}}{\partial x^\nu} \frac{\partial^2 x^\mu}{\partial x^{\bar{\beta}} \partial x^{\bar{\alpha}}} A^{\bar{\beta}}$$



# Covariant derivative

## Notation

$$\frac{\partial}{\partial x^\mu} A = \partial_\mu A = A_{,\mu}$$

- Partial derivatives of scalar fields are  $\binom{0}{1}$ -tensors

$$\phi_{,\mu} = \left( \frac{\partial x^{\bar{\alpha}}}{\partial x^\mu} \right) \phi_{,\bar{\alpha}}$$

- Partial derivatives of other tensor fields are not tensors

$$A^\mu_{,\nu} = \frac{\partial x^{\bar{\alpha}}}{\partial x^\nu} \frac{\partial}{\partial x^{\bar{\alpha}}} \left( \frac{\partial x^\mu}{\partial x^{\bar{\beta}}} A^{\bar{\beta}} \right) = \frac{\partial x^{\bar{\alpha}}}{\partial x^\nu} \frac{\partial x^\mu}{\partial x^{\bar{\beta}}} A^{\bar{\beta}}_{,\bar{\alpha}} + \frac{\partial x^{\bar{\alpha}}}{\partial x^\nu} \frac{\partial^2 x^\mu}{\partial x^{\bar{\beta}} \partial x^{\bar{\alpha}}} A^{\bar{\beta}}$$



# Covariant derivative

## Notation

$$\frac{\partial}{\partial x^\mu} A = \partial_\mu A = A_{,\mu}$$

- Partial derivatives of scalar fields are  $\binom{0}{1}$ -tensors

$$\phi_{,\mu} = \left( \frac{\partial x^{\bar{\alpha}}}{\partial x^\mu} \right) \phi_{,\bar{\alpha}}$$

- Partial derivatives of other tensor fields are not tensors

$$A^\mu_{,\nu} = \frac{\partial x^{\bar{\alpha}}}{\partial x^\nu} \frac{\partial}{\partial x^{\bar{\alpha}}} \left( \frac{\partial x^\mu}{\partial x^{\bar{\beta}}} A^{\bar{\beta}} \right) = \frac{\partial x^{\bar{\alpha}}}{\partial x^\nu} \frac{\partial x^\mu}{\partial x^{\bar{\beta}}} A^{\bar{\beta}}_{,\bar{\alpha}} + \frac{\partial x^{\bar{\alpha}}}{\partial x^\nu} \frac{\partial^2 x^\mu}{\partial x^{\bar{\beta}} \partial x^{\bar{\alpha}}} A^{\bar{\beta}}$$





## Covariant derivative

## Notation

$$\frac{\partial}{\partial x^\mu} A = \partial_\mu A = A_{,\mu}$$

- Partial derivatives of scalar fields are  $\binom{0}{1}$ -tensors

$$\phi_{,\mu} = \left( \frac{\partial x^{\bar{\alpha}}}{\partial x^\mu} \right) \phi_{,\bar{\alpha}}$$

- Partial derivatives of other tensor fields are not tensors

$$A^\mu_{,\nu} = \frac{\partial x^{\bar{\alpha}}}{\partial x^\nu} \frac{\partial}{\partial x^{\bar{\alpha}}} \left( \frac{\partial x^\mu}{\partial x^{\bar{\beta}}} A^{\bar{\beta}} \right) = \frac{\partial x^{\bar{\alpha}}}{\partial x^\nu} \frac{\partial x^\mu}{\partial x^{\bar{\beta}}} A^{\bar{\beta}}_{,\bar{\alpha}} + \frac{\partial x^{\bar{\alpha}}}{\partial x^\nu} \frac{\partial^2 x^\mu}{\partial x^{\bar{\beta}} \partial x^{\bar{\alpha}}} A^{\bar{\beta}}$$



## Covariant derivative

## Notation

$$\frac{\partial}{\partial x^\mu} A = \partial_\mu A = A_{,\mu}$$

- Partial derivatives of scalar fields are  $\binom{0}{1}$ -tensors

$$\phi_{,\mu} = \left( \frac{\partial x^{\bar{\alpha}}}{\partial x^\mu} \right) \phi_{,\bar{\alpha}}$$

- Partial derivatives of other tensor fields are not tensors

$$A^\mu_{,\nu} = \frac{\partial x^{\bar{\alpha}}}{\partial x^\nu} \frac{\partial}{\partial x^{\bar{\alpha}}} \left( \frac{\partial x^\mu}{\partial x^{\bar{\beta}}} A^{\bar{\beta}} \right) = \frac{\partial x^{\bar{\alpha}}}{\partial x^\nu} \frac{\partial x^\mu}{\partial x^{\bar{\beta}}} A^{\bar{\beta}}_{,\bar{\alpha}} + \frac{\partial x^{\bar{\alpha}}}{\partial x^\nu} \frac{\partial^2 x^\mu}{\partial x^{\bar{\beta}} \partial x^{\bar{\alpha}}} A^{\bar{\beta}}$$



# Covariant derivatives

- Partial derivatives do not transform as tensors

$$A^{\mu}{}_{,\nu} = \frac{\partial x^{\bar{\alpha}}}{\partial x^{\nu}} \frac{\partial}{\partial x^{\bar{\alpha}}} \left( \frac{\partial x^{\mu}}{\partial x^{\bar{\beta}}} A^{\bar{\beta}} \right) = \frac{\partial x^{\bar{\alpha}}}{\partial x^{\nu}} \frac{\partial x^{\mu}}{\partial x^{\bar{\beta}}} A^{\bar{\beta}}{}_{,\bar{\alpha}} + \frac{\partial x^{\bar{\alpha}}}{\partial x^{\nu}} \frac{\partial^2 x^{\mu}}{\partial x^{\bar{\beta}} \partial x^{\bar{\alpha}}} A^{\bar{\beta}}$$

- Unwanted term proportional to  $A^{\bar{\beta}}$

→ Introduce a correction factor to cancel

→ Covariant derivative,  $\nabla$  which transforms as a tensor:

$$\nabla_{\mu} A^{\nu} = \partial_{\mu} A^{\nu} + \Gamma^{\nu}_{\mu\alpha} A^{\alpha}$$

- This correction term,  $\Gamma^{\mu}_{\alpha\beta}$  is called the connection



# Covariant derivatives

- Partial derivatives do not transform as tensors

$$A^{\mu}_{,\nu} = \frac{\partial x^{\bar{\alpha}}}{\partial x^{\nu}} \frac{\partial}{\partial x^{\bar{\alpha}}} \left( \frac{\partial x^{\mu}}{\partial x^{\bar{\beta}}} A^{\bar{\beta}} \right) = \frac{\partial x^{\bar{\alpha}}}{\partial x^{\nu}} \frac{\partial x^{\mu}}{\partial x^{\bar{\beta}}} A^{\bar{\beta}}_{,\bar{\alpha}} + \frac{\partial x^{\bar{\alpha}}}{\partial x^{\nu}} \frac{\partial^2 x^{\mu}}{\partial x^{\bar{\beta}} \partial x^{\bar{\alpha}}} A^{\bar{\beta}}$$

- Unwanted term proportional to  $A^{\bar{\beta}}$

→ Introduce a correction factor to cancel

→ Covariant derivative,  $\nabla$  which transforms as a tensor:

$$\nabla_{\mu} A^{\nu} = \partial_{\mu} A^{\nu} + \Gamma^{\nu}_{\mu\alpha} A^{\alpha}$$

- This correction term,  $\Gamma^{\mu}_{\alpha\beta}$  is called the connection



# Covariant derivatives

- Partial derivatives do not transform as tensors

$$A^{\mu}{}_{,\nu} = \frac{\partial x^{\bar{\alpha}}}{\partial x^{\nu}} \frac{\partial}{\partial x^{\bar{\alpha}}} \left( \frac{\partial x^{\mu}}{\partial x^{\bar{\beta}}} A^{\bar{\beta}} \right) = \frac{\partial x^{\bar{\alpha}}}{\partial x^{\nu}} \frac{\partial x^{\mu}}{\partial x^{\bar{\beta}}} A^{\bar{\beta}}{}_{,\bar{\alpha}} + \frac{\partial x^{\bar{\alpha}}}{\partial x^{\nu}} \frac{\partial^2 x^{\mu}}{\partial x^{\bar{\beta}} \partial x^{\bar{\alpha}}} A^{\bar{\beta}}$$

- Unwanted term proportional to  $A^{\bar{\beta}}$
- Introduce a correction factor to cancel
- Covariant derivative,  $\nabla$  which transforms as a tensor:

$$\nabla_{\mu} A^{\nu} = \partial_{\mu} A^{\nu} + \Gamma^{\nu}_{\mu\alpha} A^{\alpha}$$

- This correction term,  $\Gamma^{\mu}_{\alpha\beta}$  is called the connection



# Covariant derivatives

- Partial derivatives do not transform as tensors

$$A^{\mu}{}_{,\nu} = \frac{\partial x^{\bar{\alpha}}}{\partial x^{\nu}} \frac{\partial}{\partial x^{\bar{\alpha}}} \left( \frac{\partial x^{\mu}}{\partial x^{\bar{\beta}}} A^{\bar{\beta}} \right) = \frac{\partial x^{\bar{\alpha}}}{\partial x^{\nu}} \frac{\partial x^{\mu}}{\partial x^{\bar{\beta}}} A^{\bar{\beta}}{}_{,\bar{\alpha}} + \frac{\partial x^{\bar{\alpha}}}{\partial x^{\nu}} \frac{\partial^2 x^{\mu}}{\partial x^{\bar{\beta}} \partial x^{\bar{\alpha}}} A^{\bar{\beta}}$$

- Unwanted term proportional to  $A^{\bar{\beta}}$
- Introduce a correction factor to cancel
- Covariant derivative,  $\nabla$  which transforms as a tensor:

$$\nabla_{\mu} A^{\nu} = \partial_{\mu} A^{\nu} + \Gamma^{\nu}_{\mu\alpha} A^{\alpha}$$

- This correction term,  $\Gamma^{\mu}_{\alpha\beta}$  is called the connection



# Covariant derivatives

- Partial derivatives do not transform as tensors

$$A^{\mu}{}_{,\nu} = \frac{\partial x^{\bar{\alpha}}}{\partial x^{\nu}} \frac{\partial}{\partial x^{\bar{\alpha}}} \left( \frac{\partial x^{\mu}}{\partial x^{\bar{\beta}}} A^{\bar{\beta}} \right) = \frac{\partial x^{\bar{\alpha}}}{\partial x^{\nu}} \frac{\partial x^{\mu}}{\partial x^{\bar{\beta}}} A^{\bar{\beta}}{}_{,\bar{\alpha}} + \frac{\partial x^{\bar{\alpha}}}{\partial x^{\nu}} \frac{\partial^2 x^{\mu}}{\partial x^{\bar{\beta}} \partial x^{\bar{\alpha}}} A^{\bar{\beta}}$$

- Unwanted term proportional to  $A^{\bar{\beta}}$
- Introduce a correction factor to cancel
- Covariant derivative,  $\nabla$  which transforms as a tensor:

$$\nabla_{\mu} A^{\nu} = \partial_{\mu} A^{\nu} + \Gamma^{\nu}_{\mu\alpha} A^{\alpha}$$

- This correction term,  $\Gamma^{\mu}_{\alpha\beta}$  is called the connection



# Covariant derivatives

- Partial derivatives do not transform as tensors

$$A^{\mu}{}_{,\nu} = \frac{\partial x^{\bar{\alpha}}}{\partial x^{\nu}} \frac{\partial}{\partial x^{\bar{\alpha}}} \left( \frac{\partial x^{\mu}}{\partial x^{\bar{\beta}}} A^{\bar{\beta}} \right) = \frac{\partial x^{\bar{\alpha}}}{\partial x^{\nu}} \frac{\partial x^{\mu}}{\partial x^{\bar{\beta}}} A^{\bar{\beta}}{}_{,\bar{\alpha}} + \frac{\partial x^{\bar{\alpha}}}{\partial x^{\nu}} \frac{\partial^2 x^{\mu}}{\partial x^{\bar{\beta}} \partial x^{\bar{\alpha}}} A^{\bar{\beta}}$$

- Unwanted term proportional to  $A^{\bar{\beta}}$
- Introduce a correction factor to cancel
- Covariant derivative,  $\nabla$  which transforms as a tensor:

$$\nabla_{\mu} A^{\nu} = \partial_{\mu} A^{\nu} + \Gamma^{\nu}_{\mu\alpha} A^{\alpha}$$

- This correction term,  $\Gamma^{\mu}_{\alpha\beta}$  is called the connection





# Covariant Derivative

- For the covariant derivative

$$\nabla_{\mu} A^{\nu} = \partial_{\mu} A^{\nu} + \Gamma_{\mu\alpha}^{\nu} A^{\alpha}$$

to be a  $\binom{1}{1}$ -tensor you can check the connection must transform as

$$\Gamma_{\bar{\mu}\bar{\nu}}^{\bar{\alpha}} = \frac{\partial x^{\bar{\alpha}}}{\partial x^{\rho}} \frac{\partial x^{\sigma}}{\partial x^{\bar{\mu}}} \frac{\partial x^{\lambda}}{\partial x^{\bar{\nu}}} \Gamma_{\sigma\lambda}^{\rho} - \frac{\partial x^{\rho}}{\partial x^{\bar{\mu}}} \frac{\partial x^{\sigma}}{\partial x^{\bar{\nu}}} \frac{\partial^2 x^{\bar{\alpha}}}{\partial x^{\rho} \partial x^{\sigma}}$$

- Not a tensor !
- Ugly 2nd terms cancels ugly term from partial derivative



# Covariant Derivative

- For the covariant derivative

$$\nabla_{\mu} A^{\nu} = \partial_{\mu} A^{\nu} + \Gamma_{\mu\alpha}^{\nu} A^{\alpha}$$

to be a  $\binom{1}{1}$ -tensor you can check the connection must transform as

$$\Gamma_{\bar{\mu}\bar{\nu}}^{\bar{\alpha}} = \frac{\partial x^{\bar{\alpha}}}{\partial x^{\rho}} \frac{\partial x^{\sigma}}{\partial x^{\bar{\mu}}} \frac{\partial x^{\lambda}}{\partial x^{\bar{\nu}}} \Gamma_{\sigma\lambda}^{\rho} - \frac{\partial x^{\rho}}{\partial x^{\bar{\mu}}} \frac{\partial x^{\sigma}}{\partial x^{\bar{\nu}}} \frac{\partial^2 x^{\bar{\alpha}}}{\partial x^{\rho} \partial x^{\sigma}}$$

- Not a tensor !
- Ugly 2nd terms cancels ugly term from partial derivative



# Covariant Derivative

- For the covariant derivative

$$\nabla_{\mu} A^{\nu} = \partial_{\mu} A^{\nu} + \Gamma_{\mu\alpha}^{\nu} A^{\alpha}$$

to be a  $\binom{1}{1}$ -tensor you can check the connection must transform as

$$\Gamma_{\bar{\mu}\bar{\nu}}^{\bar{\alpha}} = \frac{\partial x^{\bar{\alpha}}}{\partial x^{\rho}} \frac{\partial x^{\sigma}}{\partial x^{\bar{\mu}}} \frac{\partial x^{\lambda}}{\partial x^{\bar{\nu}}} \Gamma_{\sigma\lambda}^{\rho} - \frac{\partial x^{\rho}}{\partial x^{\bar{\mu}}} \frac{\partial x^{\sigma}}{\partial x^{\bar{\nu}}} \frac{\partial^2 x^{\bar{\alpha}}}{\partial x^{\rho} \partial x^{\sigma}}$$

- Not a tensor !
- Ugly 2nd terms cancels ugly term from partial derivative



# Covariant derivative

- For a,  $\binom{0}{1}$  tensor,  $\omega_\alpha$ , we can also make a covariant derivative

$$\nabla_\nu \omega_\mu = \partial_\nu \omega_\mu - \Gamma_{\nu\mu}^\alpha \omega_\alpha$$

- For a general  $\binom{n}{m}$  tensor the rule is

$$\begin{aligned} \nabla_\alpha T^{\mu_1 \mu_2 \dots \mu_n}_{\nu_1 \dots \nu_m} &= \partial_\alpha T^{\mu_1 \mu_2 \dots \mu_n}_{\nu_1 \dots \nu_m} + \Gamma_{\alpha\lambda}^{\mu_1} T^{\lambda \mu_2 \dots \mu_n}_{\nu_1 \dots \nu_m} \\ &\quad + \Gamma_{\alpha\lambda}^{\mu_2} T^{\mu_1 \lambda \dots \mu_n}_{\nu_1 \dots \nu_m} + \dots \\ &\quad - \Gamma_{\alpha\nu_1}^\lambda T^{\mu_1 \mu_2 \dots \mu_n}_{\lambda \dots \nu_m} \dots \end{aligned}$$

- Rule:  $+\Gamma$  for each upper index and  $-\Gamma$  for each lower index.



# Covariant derivative

- For a,  $\binom{0}{1}$  tensor,  $\omega_\alpha$ , we can also make a covariant derivative

$$\nabla_\nu \omega_\mu = \partial_\nu \omega_\mu - \Gamma_{\nu\mu}^\alpha \omega_\alpha$$

- For a general  $\binom{n}{m}$  tensor the rule is

$$\begin{aligned} \nabla_\alpha T^{\mu_1 \mu_2 \dots \mu_n}_{\nu_1 \dots \nu_m} &= \partial_\alpha T^{\mu_1 \mu_2 \dots \mu_n}_{\nu_1 \dots \nu_m} + \Gamma_{\alpha\lambda}^{\mu_1} T^{\lambda \mu_2 \dots \mu_n}_{\nu_1 \dots \nu_m} \\ &\quad + \Gamma_{\alpha\lambda}^{\mu_2} T^{\mu_1 \lambda \dots \mu_n}_{\nu_1 \dots \nu_m} + \dots \\ &\quad - \Gamma_{\alpha\nu_1}^\lambda T^{\mu_1 \mu_2 \dots \mu_n}_{\lambda \dots \nu_m} \dots \end{aligned}$$

- Rule:  $+\Gamma$  for each upper index and  $-\Gamma$  for each lower index.



# Covariant derivative

- For a,  $\binom{0}{1}$  tensor,  $\omega_\alpha$ , we can also make a covariant derivative

$$\nabla_\nu \omega_\mu = \partial_\nu \omega_\mu - \Gamma_{\nu\mu}^\alpha \omega_\alpha$$

- For a general  $\binom{n}{m}$  tensor the rule is

$$\begin{aligned} \nabla_\alpha T^{\mu_1 \mu_2 \dots \mu_n}_{\nu_1 \dots \nu_m} &= \partial_\alpha T^{\mu_1 \mu_2 \dots \mu_n}_{\nu_1 \dots \nu_m} + \Gamma_{\alpha\lambda}^{\mu_1} T^{\lambda \mu_2 \dots \mu_n}_{\nu_1 \dots \nu_m} \\ &\quad + \Gamma_{\alpha\lambda}^{\mu_2} T^{\mu_1 \lambda \dots \mu_n}_{\nu_1 \dots \nu_m} + \dots \\ &\quad - \Gamma_{\alpha\nu_1}^\lambda T^{\mu_1 \mu_2 \dots \mu_n}_{\lambda \dots \nu_m} \dots \end{aligned}$$

- Rule:  $+\Gamma$  for each upper index and  $-\Gamma$  for each lower index.



# Christoffel Symbols

- What are these mysterious  $\Gamma$ 's?
- Many choices
- Special choice = Christoffel Symbols

$$\Gamma_{\alpha\beta}^{\gamma} = \frac{1}{2}g^{\gamma\mu}(g_{\mu\alpha,\beta} + g_{\mu\beta,\alpha} - g_{\alpha\beta,\mu})$$

- Can check it transforms in the right way
- Also satisfies:  $\nabla_{\alpha}g_{\mu\nu} = 0$     $\nabla_{\alpha}g^{\mu\nu} = 0$
- Notice: In flat Cartesian coordinates  $\Gamma_{\mu\nu}^{\alpha} = 0$  so

$$\nabla_{\mu} = \partial_{\mu}$$

in that case.

- Will give some justification for Christoffels later



# Christoffel Symbols

- What are these mysterious  $\Gamma$ 's?
- Many choices
- Special choice = Christoffel Symbols

$$\Gamma_{\alpha\beta}^{\gamma} = \frac{1}{2}g^{\gamma\mu}(g_{\mu\alpha,\beta} + g_{\mu\beta,\alpha} - g_{\alpha\beta,\mu})$$

- Can check it transforms in the right way
- Also satisfies:  $\nabla_{\alpha}g_{\mu\nu} = 0$     $\nabla_{\alpha}g^{\mu\nu} = 0$
- Notice: In flat Cartesian coordinates  $\Gamma_{\mu\nu}^{\alpha} = 0$  so

$$\nabla_{\mu} = \partial_{\mu}$$

in that case.

- Will give some justification for Christoffels later





# Christoffel Symbols

- What are these mysterious  $\Gamma$ 's?
- Many choices
- Special choice = Christoffel Symbols

$$\Gamma_{\alpha\beta}^{\gamma} = \frac{1}{2}g^{\gamma\mu}(g_{\mu\alpha,\beta} + g_{\mu\beta,\alpha} - g_{\alpha\beta,\mu})$$

- Can check it transforms in the right way
- Also satisfies:  $\nabla_{\alpha}g_{\mu\nu} = 0$     $\nabla_{\alpha}g^{\mu\nu} = 0$
- Notice: In flat Cartesian coordinates  $\Gamma_{\mu\nu}^{\alpha} = 0$  so

$$\nabla_{\mu} = \partial_{\mu}$$

in that case.

- Will give some justification for Christoffels later



# Christoffel Symbols

- What are these mysterious  $\Gamma$ 's?
- Many choices
- Special choice = Christoffel Symbols

$$\Gamma_{\alpha\beta}^{\gamma} = \frac{1}{2}g^{\gamma\mu}(g_{\mu\alpha,\beta} + g_{\mu\beta,\alpha} - g_{\alpha\beta,\mu})$$

- Can check it transforms in the right way
- Also satisfies:  $\nabla_{\alpha}g_{\mu\nu} = 0$   $\nabla_{\alpha}g^{\mu\nu} = 0$
- Notice: In flat Cartesian coordinates  $\Gamma_{\mu\nu}^{\alpha} = 0$  so

$$\nabla_{\mu} = \partial_{\mu}$$

in that case.

- Will give some justification for Christoffels later



# Christoffel Symbols

- What are these mysterious  $\Gamma$ 's?
- Many choices
- Special choice = Christoffel Symbols

$$\Gamma_{\alpha\beta}^{\gamma} = \frac{1}{2}g^{\gamma\mu}(g_{\mu\alpha,\beta} + g_{\mu\beta,\alpha} - g_{\alpha\beta,\mu})$$

- Can check it transforms in the right way
- Also satisfies:  $\nabla_{\alpha}g_{\mu\nu} = 0$   $\nabla_{\alpha}g^{\mu\nu} = 0$
- Notice: In flat Cartesian coordinates  $\Gamma_{\mu\nu}^{\alpha} = 0$  so

$$\nabla_{\mu} = \partial_{\mu}$$

in that case.

- Will give some justification for Christoffels later



# Christoffel Symbols

- What are these mysterious  $\Gamma$ 's?
- Many choices
- Special choice = Christoffel Symbols

$$\Gamma_{\alpha\beta}^{\gamma} = \frac{1}{2}g^{\gamma\mu}(g_{\mu\alpha,\beta} + g_{\mu\beta,\alpha} - g_{\alpha\beta,\mu})$$

- Can check it transforms in the right way
- Also satisfies:  $\nabla_{\alpha}g_{\mu\nu} = 0$     $\nabla_{\alpha}g^{\mu\nu} = 0$
- Notice: In flat Cartesian coordinates  $\Gamma_{\mu\nu}^{\alpha} = 0$  so

$$\nabla_{\mu} = \partial_{\mu}$$

in that case.

- Will give some justification for Christoffels later



# Christoffel Symbols

- What are these mysterious  $\Gamma$ 's?
- Many choices
- Special choice = Christoffel Symbols

$$\Gamma_{\alpha\beta}^{\gamma} = \frac{1}{2}g^{\gamma\mu}(g_{\mu\alpha,\beta} + g_{\mu\beta,\alpha} - g_{\alpha\beta,\mu})$$

- Can check it transforms in the right way
- Also satisfies:  $\nabla_{\alpha}g_{\mu\nu} = 0$     $\nabla_{\alpha}g^{\mu\nu} = 0$
- Notice: In flat Cartesian coordinates  $\Gamma_{\mu\nu}^{\alpha} = 0$  so

$$\nabla_{\mu} = \partial_{\mu}$$

in that case.

- Will give some justification for Christoffels later



# Christoffel Symbols

- What are these mysterious  $\Gamma$ 's?
- Many choices
- Special choice = Christoffel Symbols

$$\Gamma_{\alpha\beta}^{\gamma} = \frac{1}{2}g^{\gamma\mu}(g_{\mu\alpha,\beta} + g_{\mu\beta,\alpha} - g_{\alpha\beta,\mu})$$

- Can check it transforms in the right way
- Also satisfies:  $\nabla_{\alpha}g_{\mu\nu} = 0$     $\nabla_{\alpha}g^{\mu\nu} = 0$
- Notice: In flat Cartesian coordinates  $\Gamma_{\mu\nu}^{\alpha} = 0$  so

$$\nabla_{\mu} = \partial_{\mu}$$

in that case.

- Will give some justification for Christoffels later



# Christoffel Symbols

- What are these mysterious  $\Gamma$ 's?
- Many choices
- Special choice = Christoffel Symbols

$$\Gamma_{\alpha\beta}^{\gamma} = \frac{1}{2}g^{\gamma\mu}(g_{\mu\alpha,\beta} + g_{\mu\beta,\alpha} - g_{\alpha\beta,\mu})$$

- Can check it transforms in the right way
- Also satisfies:  $\nabla_{\alpha}g_{\mu\nu} = 0$     $\nabla_{\alpha}g^{\mu\nu} = 0$
- Notice: In flat Cartesian coordinates  $\Gamma_{\mu\nu}^{\alpha} = 0$  so

$$\nabla_{\mu} = \partial_{\mu}$$

in that case.

- Will give some justification for Christoffels later



# Riemann Curvature Tensor

- Ordinary derivatives commute :  $[\partial_\alpha, \partial_\beta] = 0$
- Failure of  $\nabla$ 's to commute related to failure of E5
- Commutator of  $\nabla$ 's encodes curvature.
- Riemann Curvature tensor

$$[\nabla_\alpha, \nabla_\beta]X^\gamma = R^\gamma_{\lambda\alpha\beta}X^\lambda$$

$\Rightarrow$

$$R^\gamma_{\lambda\alpha\beta} = \partial_\alpha \Gamma^\gamma_{\beta\lambda} - \partial_\beta \Gamma^\gamma_{\alpha\lambda} + \Gamma^\gamma_{\alpha\rho} \Gamma^\rho_{\beta\lambda} - \Gamma^\gamma_{\beta\rho} \Gamma^\rho_{\alpha\lambda}$$





# Riemann Curvature Tensor

- Ordinary derivatives commute :  $[\partial_\alpha, \partial_\beta] = 0$
- Failure of  $\nabla$ 's to commute related to failure of E5
- Commutator of  $\nabla$ 's encodes curvature.
- Riemann Curvature tensor

$$[\nabla_\alpha, \nabla_\beta]X^\gamma = R^\gamma_{\lambda\alpha\beta}X^\lambda$$

$\Rightarrow$

$$R^\gamma_{\lambda\alpha\beta} = \partial_\alpha \Gamma^\gamma_{\beta\lambda} - \partial_\beta \Gamma^\gamma_{\alpha\lambda} + \Gamma^\gamma_{\alpha\rho} \Gamma^\rho_{\beta\lambda} - \Gamma^\gamma_{\beta\rho} \Gamma^\rho_{\alpha\lambda}$$



# Riemann Curvature Tensor

- Ordinary derivatives commute :  $[\partial_\alpha, \partial_\beta] = 0$
- Failure of  $\nabla$ 's to commute related to failure of E5
- Commutator of  $\nabla$ 's encodes curvature.
- Riemann Curvature tensor

$$[\nabla_\alpha, \nabla_\beta]X^\gamma = R^\gamma_{\lambda\alpha\beta}X^\lambda$$

$\Rightarrow$

$$R^\gamma_{\lambda\alpha\beta} = \partial_\alpha \Gamma^\gamma_{\beta\lambda} - \partial_\beta \Gamma^\gamma_{\alpha\lambda} + \Gamma^\gamma_{\alpha\rho} \Gamma^\rho_{\beta\lambda} - \Gamma^\gamma_{\beta\rho} \Gamma^\rho_{\alpha\lambda}$$



# Riemann Curvature Tensor

- Ordinary derivatives commute :  $[\partial_\alpha, \partial_\beta] = 0$
- Failure of  $\nabla$ 's to commute related to failure of E5
- Commutator of  $\nabla$ 's encodes curvature.
- Riemann Curvature tensor

$$[\nabla_\alpha, \nabla_\beta]X^\gamma = R^\gamma_{\lambda\alpha\beta}X^\lambda$$

$\Rightarrow$

$$R^\gamma_{\lambda\alpha\beta} = \partial_\alpha \Gamma^\gamma_{\beta\lambda} - \partial_\beta \Gamma^\gamma_{\alpha\lambda} + \Gamma^\gamma_{\alpha\rho} \Gamma^\rho_{\beta\lambda} - \Gamma^\gamma_{\beta\rho} \Gamma^\rho_{\alpha\lambda}$$



# Riemann Curvature Tensor

- Ordinary derivatives commute :  $[\partial_\alpha, \partial_\beta] = 0$
- Failure of  $\nabla$ 's to commute related to failure of E5
- Commutator of  $\nabla$ 's encodes curvature.
- Riemann Curvature tensor

$$[\nabla_\alpha, \nabla_\beta]X^\gamma = R^\gamma_{\lambda\alpha\beta}X^\lambda$$

$\Rightarrow$

$$R^\gamma_{\lambda\alpha\beta} = \partial_\alpha\Gamma^\gamma_{\beta\lambda} - \partial_\beta\Gamma^\gamma_{\alpha\lambda} + \Gamma^\gamma_{\alpha\rho}\Gamma^\rho_{\beta\lambda} - \Gamma^\gamma_{\beta\rho}\Gamma^\rho_{\alpha\lambda}$$



# Notation: Ricci Tensor

- The trace of the Riemann Tensor gives us Ricci Tensor

$$R_{\alpha\beta} = \sum_{\gamma=0\dots3} R^{\gamma}{}_{\alpha\gamma\beta}$$

- The trace of the Ricci Tensor gives us the Ricci Scalar

$$R = R^{\lambda}{}_{\lambda} = g^{\lambda\nu} R_{\nu\lambda}$$



# Notation: Ricci Tensor

- The trace of the Riemann Tensor gives us Ricci Tensor

$$R_{\alpha\beta} = R^{\gamma}_{\alpha\gamma\beta}$$

- The trace of the Ricci Tensor gives us the Ricci Scalar

$$R = R^{\lambda}_{\lambda} = g^{\lambda\nu} R_{\nu\lambda}$$



# Notation: Ricci Tensor

- The trace of the Riemann Tensor gives us Ricci Tensor

$$R_{\alpha\beta} = R^{\gamma}_{\alpha\gamma\beta}$$

- The trace of the Ricci Tensor gives us the Ricci Scalar

$$R = R^{\lambda}_{\lambda} = g^{\lambda\nu} R_{\nu\lambda}$$



- The Riemann tensor has many symmetries and satisfies many identities (which follow from the definition)
- NB is the Bianchi identity

$$\nabla_{\lambda} R_{\alpha\beta\mu\nu} + \nabla_{\alpha} R_{\beta\lambda\mu\nu} + \nabla_{\beta} R_{\lambda\alpha\mu\nu} = 0$$

- This implies that the Einstein Tensor

$$G^{\mu\nu} = R^{\mu\nu} - \frac{1}{2}g^{\mu\nu}R$$

satisfies

- $\nabla_{\mu} G^{\mu\nu} = 0$





- The Riemann tensor has many symmetries and satisfies many identities (which follow from the definition)
- NB is the Bianchi identity

$$\nabla_{\lambda} R_{\alpha\beta\mu\nu} + \nabla_{\alpha} R_{\beta\lambda\mu\nu} + \nabla_{\beta} R_{\lambda\alpha\mu\nu} = 0$$

- This implies that the Einstein Tensor

$$G^{\mu\nu} = R^{\mu\nu} - \frac{1}{2}g^{\mu\nu}R$$

satisfies

- $\nabla_{\mu} G^{\mu\nu} = 0$



- The Riemann tensor has many symmetries and satisfies many identities (which follow from the definition)
- NB is the Bianchi identity

$$\nabla_{\lambda} R_{\alpha\beta\mu\nu} + \nabla_{\alpha} R_{\beta\lambda\mu\nu} + \nabla_{\beta} R_{\lambda\alpha\mu\nu} = 0$$

- This implies that the Einstein Tensor

$$G^{\mu\nu} = R^{\mu\nu} - \frac{1}{2}g^{\mu\nu}R$$

satisfies

- $\nabla_{\mu} G^{\mu\nu} = 0$



- The Riemann tensor has many symmetries and satisfies many identities (which follow from the definition)
- NB is the Bianchi identity

$$\nabla_{\lambda} R_{\alpha\beta\mu\nu} + \nabla_{\alpha} R_{\beta\lambda\mu\nu} + \nabla_{\beta} R_{\lambda\alpha\mu\nu} = 0$$

- This implies that the Einstein Tensor

$$G^{\mu\nu} = R^{\mu\nu} - \frac{1}{2}g^{\mu\nu}R$$

satisfies

- $\nabla_{\mu} G^{\mu\nu} = 0$



# “Straight lines” in curved space

- How can we generalise the concept of a straight line to curved space?

- Straight line = shortest distance two points

→ Extremise the distance (or proper time)

$$= \int |ds| = \int d\lambda \sqrt{\left| g_{\mu\nu} \frac{dx^\mu}{d\lambda} \frac{dx^\nu}{d\lambda} \right|}$$

⇒ Geodesic equation

$$\frac{d^2 x^\mu}{d\lambda^2} + \Gamma_{\alpha\beta}^\mu \frac{dx^\alpha}{d\lambda} \frac{dx^\beta}{d\lambda} = 0$$

- $\Gamma_{\rho\sigma}^\alpha$  = Christoffel symbols
- $\lambda$  must be linearly related proper-distance/proper-time.



# “Straight lines” in curved space

- How can we generalise the concept of a straight line to curved space?
  - Straight line = shortest distance two points

→ Extremise the distance (or proper time)

$$= \int |ds| = \int d\lambda \sqrt{|g_{\mu\nu} \frac{dx^\mu}{d\lambda} \frac{dx^\nu}{d\lambda}|}$$

⇒ Geodesic equation

$$\frac{d^2 x^\mu}{d\lambda^2} + \Gamma_{\alpha\beta}^\mu \frac{dx^\alpha}{d\lambda} \frac{dx^\beta}{d\lambda} = 0$$

- $\Gamma_{\rho\sigma}^\alpha$  = Christoffel symbols
- $\lambda$  must be linearly related proper-distance/proper-time.



# “Straight lines” in curved space

- How can we generalise the concept of a straight line to curved space?
  - Straight line = shortest distance two points
- Extremise the distance (or proper time)

$$= \int |ds| = \int d\lambda \sqrt{|g_{\mu\nu} \frac{dx^\mu}{d\lambda} \frac{dx^\nu}{d\lambda}|}$$

⇒ Geodesic equation

$$\frac{d^2 x^\mu}{d\lambda^2} + \Gamma_{\alpha\beta}^\mu \frac{dx^\alpha}{d\lambda} \frac{dx^\beta}{d\lambda} = 0$$

- $\Gamma_{\rho\sigma}^\alpha$  = Christoffel symbols
- $\lambda$  must be linearly related proper-distance/proper-time.



# “Straight lines” in curved space

- How can we generalise the concept of a straight line to curved space?
  - Straight line = shortest distance two points
- Extremise the distance (or proper time)

$$= \int |ds| = \int d\lambda \sqrt{\left| g_{\mu\nu} \frac{dx^\mu}{d\lambda} \frac{dx^\nu}{d\lambda} \right|}$$

⇒ Geodesic equation

$$\frac{d^2 x^\mu}{d\lambda^2} + \Gamma_{\alpha\beta}^\mu \frac{dx^\alpha}{d\lambda} \frac{dx^\beta}{d\lambda} = 0$$

- $\Gamma_{\rho\sigma}^\alpha$  = Christoffel symbols
- $\lambda$  must be linearly related proper-distance/proper-time.



# “Straight lines” in curved space

- How can we generalise the concept of a straight line to curved space?
  - Straight line = shortest distance two points
- Extremise the distance (or proper time)
- =  $\int |ds| = \int d\lambda \sqrt{|g_{\mu\nu} \frac{dx^\mu}{d\lambda} \frac{dx^\nu}{d\lambda}|}$
- ⇒ Geodesic equation

$$\frac{d^2 x^\mu}{d\lambda^2} + \Gamma_{\alpha\beta}^\mu \frac{dx^\alpha}{d\lambda} \frac{dx^\beta}{d\lambda} = 0$$

- $\Gamma_{\rho\sigma}^\alpha$  = Christoffel symbols
- $\lambda$  must be linearly related proper-distance/proper-time.





# “Straight lines” in curved space

- How can we generalise the concept of a straight line to curved space?
  - Straight line = shortest distance two points
- Extremise the distance (or proper time)
- =  $\int |ds| = \int d\lambda \sqrt{|g_{\mu\nu} \frac{dx^\mu}{d\lambda} \frac{dx^\nu}{d\lambda}|}$
- ⇒ Geodesic equation

$$\frac{d^2 x^\mu}{d\lambda^2} + \Gamma_{\alpha\beta}^\mu \frac{dx^\alpha}{d\lambda} \frac{dx^\beta}{d\lambda} = 0$$

- $\Gamma_{\rho\sigma}^\alpha$  = Christoffel symbols
- $\lambda$  must be linearly related proper-distance/proper-time.



# “Straight lines” in curved space

- How can we generalise the concept of a straight line to curved space?
  - Straight line = shortest distance two points
- Extremise the distance (or proper time)
- =  $\int |ds| = \int d\lambda \sqrt{|g_{\mu\nu} \frac{dx^\mu}{d\lambda} \frac{dx^\nu}{d\lambda}|}$
- ⇒ Geodesic equation

$$\frac{d^2 x^\mu}{d\lambda^2} + \Gamma_{\alpha\beta}^\mu \frac{dx^\alpha}{d\lambda} \frac{dx^\beta}{d\lambda} = 0$$

- $\Gamma_{\rho\sigma}^\alpha$  = Christoffel symbols
- $\lambda$  must be linearly related proper-distance/proper-time.



# “Straight lines”

- In flat cartesian coordinates,  $\Gamma_{\beta\gamma}^{\alpha} = 0$

→  $\frac{d^2 x^{\mu}}{d\lambda^2} = 0$

→ Straight-lines

- 

- Geodesic equation

$$\ddot{x}^{\mu} = -\Gamma_{\alpha\beta}^{\mu} \dot{x}^{\alpha} \dot{x}^{\beta}$$

- Christoffel symbols  $\leftrightarrow$  Gravitational potential
- Works because inertial mass = gravitational mass



# “Straight lines”

- In flat cartesian coordinates,  $\Gamma_{\beta\gamma}^{\alpha} = 0$

→  $\frac{d^2 x^{\mu}}{d\lambda^2} = 0$

→ Straight-lines

- 

- Geodesic equation

$$\ddot{x}^{\mu} = -\Gamma_{\alpha\beta}^{\mu} \dot{x}^{\alpha} \dot{x}^{\beta}$$

- Christoffel symbols  $\leftrightarrow$  Gravitational potential
- Works because inertial mass = gravitational mass



# “Straight lines”

- In flat cartesian coordinates,  $\Gamma_{\beta\gamma}^{\alpha} = 0$

→  $\frac{d^2 x^{\mu}}{d\lambda^2} = 0$

→ Straight-lines

- 

- Geodesic equation

$$\ddot{x}^{\mu} = -\Gamma_{\alpha\beta}^{\mu} \dot{x}^{\alpha} \dot{x}^{\beta}$$

- Christoffel symbols  $\leftrightarrow$  Gravitational potential
- Works because inertial mass = gravitational mass



# “Straight lines”

- In flat cartesian coordinates,  $\Gamma_{\beta\gamma}^{\alpha} = 0$

→  $\frac{d^2 x^{\mu}}{d\lambda^2} = 0$

→ Straight-lines

- Newton's equation

$$m\ddot{x}^i = F^i$$

- Geodesic equation

$$\ddot{x}^{\mu} = -\Gamma_{\alpha\beta}^{\mu} \dot{x}^{\alpha} \dot{x}^{\beta}$$

- Christoffel symbols  $\leftrightarrow$  Gravitational potential
- Works because inertial mass = gravitational mass



# “Straight lines”

- In flat cartesian coordinates,  $\Gamma_{\beta\gamma}^{\alpha} = 0$
- $\frac{d^2 x^{\mu}}{d\lambda^2} = 0$
- Straight-lines
- Newton’s equation for gravity

$$m\ddot{x}^i = -m\partial^i(\phi_G)$$

- Geodesic equation

$$\ddot{x}^{\mu} = -\Gamma_{\alpha\beta}^{\mu} \dot{x}^{\alpha} \dot{x}^{\beta}$$

- Christoffel symbols  $\leftrightarrow$  Gravitational potential
- Works because inertial mass = gravitational mass



# “Straight lines”

- In flat cartesian coordinates,  $\Gamma_{\beta\gamma}^{\alpha} = 0$
- $\frac{d^2 x^{\mu}}{d\lambda^2} = 0$
- Straight-lines
- Newton's equation for gravity

$$\ddot{x}^i = -\partial^i(\phi_G)$$

- Geodesic equation

$$\ddot{x}^{\mu} = -\Gamma_{\alpha\beta}^{\mu} \dot{x}^{\alpha} \dot{x}^{\beta}$$

- Christoffel symbols  $\leftrightarrow$  Gravitational potential
- Works because inertial mass = gravitational mass





# “Straight lines”

- In flat cartesian coordinates,  $\Gamma_{\beta\gamma}^{\alpha} = 0$
- $\frac{d^2 x^{\mu}}{d\lambda^2} = 0$
- Straight-lines
- Newton’s equation for gravity

$$\ddot{x}^i = -\partial^i(\phi_G)$$

- Geodesic equation

$$\ddot{x}^{\mu} = -\Gamma_{\alpha\beta}^{\mu} \dot{x}^{\alpha} \dot{x}^{\beta}$$

- Christoffel symbols  $\leftrightarrow$  Gravitational potential
- Works because inertial mass = gravitational mass



# “Straight lines”

- In flat cartesian coordinates,  $\Gamma_{\beta\gamma}^{\alpha} = 0$
- $\frac{d^2 x^{\mu}}{d\lambda^2} = 0$
- Straight-lines
- Newton's equation for gravity

$$\ddot{x}^i = -\partial^i(\phi_G)$$

- Geodesic equation

$$\ddot{x}^{\mu} = -\Gamma_{\alpha\beta}^{\mu} \dot{x}^{\alpha} \dot{x}^{\beta}$$

- Christoffel symbols  $\leftrightarrow$  Gravitational potential
- Works because inertial mass = gravitational mass



# “Straight lines”

- In flat cartesian coordinates,  $\Gamma_{\beta\gamma}^{\alpha} = 0$
- $\frac{d^2 x^{\mu}}{d\lambda^2} = 0$
- Straight-lines
- Newton's equation for gravity

$$\ddot{x}^i = -\partial^i(\phi_G)$$

- Geodesic equation

$$\ddot{x}^{\mu} = -\Gamma_{\alpha\beta}^{\mu} \dot{x}^{\alpha} \dot{x}^{\beta}$$

- Christoffel symbols  $\leftrightarrow$  Gravitational potential
- Works because inertial mass = gravitational mass



# Gravity is not a force

- One of Einstein's great ideas is that gravity is not a force.
- No force  $\rightarrow$  particles move on straight lines
- In curved space, straight lines  $\rightarrow$  geodesics
- Gravity = curvature of space-time.
- No force + curved space-time  $\rightarrow$  particles move geodesics
- Hard to reconcile with QFT. Wrong?



# Gravity is not a force

- One of Einstein's great ideas is that gravity is not a force.
- No force  $\rightarrow$  particles move on straight lines
- In curved space, straight lines  $\rightarrow$  geodesics
- Gravity = curvature of space-time.
- No force + curved space-time  $\rightarrow$  particles move geodesics
- Hard to reconcile with QFT. Wrong?



# Gravity is not a force

- One of Einstein's great ideas is that gravity is not a force.
- No force  $\rightarrow$  particles move on straight lines
- In curved space, straight lines  $\rightarrow$  geodesics
- Gravity = curvature of space-time.
- No force + curved space-time  $\rightarrow$  particles move geodesics
- Hard to reconcile with QFT. Wrong?



# Gravity is not a force

- One of Einstein's great ideas is that gravity is not a force.
- No force  $\rightarrow$  particles move on straight lines
- In curved space, straight lines  $\rightarrow$  geodesics
- Gravity = curvature of space-time.
- No force + curved space-time  $\rightarrow$  particles move geodesics
- Hard to reconcile with QFT. Wrong?



# Gravity is not a force

- One of Einstein's great ideas is that gravity is not a force.
- No force  $\rightarrow$  particles move on straight lines
- In curved space, straight lines  $\rightarrow$  geodesics
- Gravity = curvature of space-time.
- No force + curved space-time  $\rightarrow$  particles move geodesics
- Hard to reconcile with QFT. Wrong?





# Gravity is not a force

- One of Einstein's great ideas is that gravity is not a force.
- No force  $\rightarrow$  particles move on straight lines
- In curved space, straight lines  $\rightarrow$  geodesics
- Gravity = curvature of space-time.
- No force + curved space-time  $\rightarrow$  particles move geodesics
- Hard to reconcile with QFT. Wrong?



# Mass-Energy curves space-time

- In Newton's theory mass is source of gravitational force
- Einstein: Gravity = curvature
- Mass-Energy curves space-time
- How? Need tensor description of matter



# Mass-Energy curves space-time

- In Newton's theory mass is source of gravitational force
  - Einstein: Gravity = curvature
- Mass-Energy curves space-time
- How? Need tensor description of matter



# Mass-Energy curves space-time

- In Newton's theory mass is source of gravitational force
- Einstein: Gravity = curvature
- Mass-Energy curves space-time
  - How? Need tensor description of matter



# Mass-Energy curves space-time

- In Newton's theory mass is source of gravitational force
- Einstein: Gravity = curvature
- Mass-Energy curves space-time
- How? Need tensor description of matter



# Mass-Energy curves space-time

- In Newton's theory mass is source of gravitational force
- Einstein: Gravity = curvature
- Mass-Energy curves space-time
- How? Need tensor description of matter



# Stress-Energy Tensor

## Stress-Energy Tensor definition

$T^{\mu\nu}$  = Flux of  $\mu^{\text{th}}$  component of momentum in  $\nu^{\text{th}}$  direction

- For example, in the rest frame of a perfect fluid in flat space-time,

$$T^{\mu\nu} = \begin{pmatrix} \rho & 0 & 0 & 0 \\ 0 & p & 0 & 0 \\ 0 & 0 & p & 0 \\ 0 & 0 & 0 & p \end{pmatrix}$$

- Conservation of mass-energy  $\Leftrightarrow \nabla_{\mu} T^{\mu\nu} = 0$



# Stress-Energy Tensor

## Stress-Energy Tensor definition

$T^{\mu\nu}$  = Flux of  $\mu^{\text{th}}$  component of momentum in  $\nu^{\text{th}}$  direction

- For example, in the rest frame of a perfect fluid in flat space-time,

$$T^{\mu\nu} = \begin{pmatrix} \rho & 0 & 0 & 0 \\ 0 & p & 0 & 0 \\ 0 & 0 & p & 0 \\ 0 & 0 & 0 & p \end{pmatrix}$$

- Conservation of mass-energy  $\Leftrightarrow \nabla_{\mu} T^{\mu\nu} = 0$





# Stress-Energy Tensor

## Stress-Energy Tensor definition

$T^{\mu\nu}$  = Flux of  $\mu^{\text{th}}$  component of momentum in  $\nu^{\text{th}}$  direction

- For example, in the rest frame of a perfect fluid in flat space-time,

$$T^{\mu\nu} = \begin{pmatrix} \rho & 0 & 0 & 0 \\ 0 & p & 0 & 0 \\ 0 & 0 & p & 0 \\ 0 & 0 & 0 & p \end{pmatrix}$$

- Conservation of mass-energy  $\Leftrightarrow \nabla_{\mu} T^{\mu\nu} = 0$



- Mass energy curves space-time
- $\nabla_{\mu} G^{\mu\nu} = 0$
- $\nabla_{\mu} T^{\mu\nu} = 0$
- Einstein guessed  $G^{\mu\nu} \propto T^{\mu\nu}$
- To recover Newtonian gravity need,

$$G^{\mu\nu} = 8\pi G T^{\mu\nu}$$



- Mass energy curves space-time
- $\nabla_{\mu} G^{\mu\nu} = 0$
- $\nabla_{\mu} T^{\mu\nu} = 0$
- Einstein guessed  $G^{\mu\nu} \propto T^{\mu\nu}$
- To recover Newtonian gravity need,

$$G^{\mu\nu} = 8\pi G T^{\mu\nu}$$



- Mass energy curves space-time
- $\nabla_{\mu} G^{\mu\nu} = 0$
- $\nabla_{\mu} T^{\mu\nu} = 0$
- Einstein guessed  $G^{\mu\nu} \propto T^{\mu\nu}$
- To recover Newtonian gravity need,

$$G^{\mu\nu} = 8\pi G T^{\mu\nu}$$



- Mass energy curves space-time
- $\nabla_{\mu} G^{\mu\nu} = 0$
- $\nabla_{\mu} T^{\mu\nu} = 0$
- Einstein guessed  $G^{\mu\nu} \propto T^{\mu\nu}$
- To recover Newtonian gravity need,

$$G^{\mu\nu} = 8\pi G T^{\mu\nu}$$



- Mass energy curves space-time
- $\nabla_{\mu} G^{\mu\nu} = 0$
- $\nabla_{\mu} T^{\mu\nu} = 0$
- Einstein guessed  $G^{\mu\nu} \propto T^{\mu\nu}$
- To recover Newtonian gravity need,

$$G^{\mu\nu} = 8\pi G T^{\mu\nu}$$



# General Relativity

## GR Summary - 3 statements:

- Space-time is a curved pseudo-Riemannian manifold
- Matter curves space-time according to Einstein's equation:

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 8\pi GT_{\mu\nu}$$

- Free particles follow geodesics in space-time.
- We need to a bit of mathematics to appreciate the summary
- Work in units where  $c = 1$



# General Relativity

## GR Summary - 3 statements:

- Space-time is a curved pseudo-Riemannian manifold
- Matter curves space-time according to Einstein's equation:

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 8\pi GT_{\mu\nu}$$

- Free particles follow geodesics in space-time.
- We need to a bit of mathematics to appreciate the summary
- Work in units where  $c = 1$





# General Relativity

## GR Summary - 3 statements:

- Space-time is a curved pseudo-Riemannian manifold
- Matter curves space-time according to Einstein's equation:

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 8\pi GT_{\mu\nu}$$

- Free particles follow geodesics in space-time.
- We need to a bit of mathematics to appreciate the summary
- Work in units where  $c = 1$



# General Relativity

## GR Summary - 3 statements:

- Space-time is a curved pseudo-Riemannian manifold
- Matter curves space-time according to Einstein's equation:

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 8\pi GT_{\mu\nu}$$

- Free particles follow geodesics in space-time.
- We need to a bit of mathematics to appreciate the summary
- Work in units where  $c = 1$



# General Relativity

## GR Summary - 3 statements:

- Space-time is a curved pseudo-Riemannian manifold
- Matter curves space-time according to Einstein's equation:

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 8\pi GT_{\mu\nu}$$

- Free particles follow geodesics in space-time.
- We need to a bit of mathematics to appreciate the summary
- Work in units where  $c = 1$



# General Relativity

## GR Summary - 3 statements:

- Space-time is a curved pseudo-Riemannian manifold
- Matter curves space-time according to Einstein's equation:

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 8\pi GT_{\mu\nu}$$

- Free particles follow geodesics in space-time.
- We need to a bit of mathematics to appreciate the summary
- Work in units where  $c = 1$





# Bibliography



S. Carrol

A No-Nonsense Introduction to General Relativity<sup>1</sup>

*<http://preposterousuniverse.com/grnotes/grtinypdf.pdf>*

---

<sup>1</sup>Lecture loosely base on these notes.

